

Continuous-variable quantum information, multi-mode quantum optics and bosonic error correcting codes

Francesco Arzani

frarzani.github.io



Alexander von Humboldt
Stiftung/Foundation

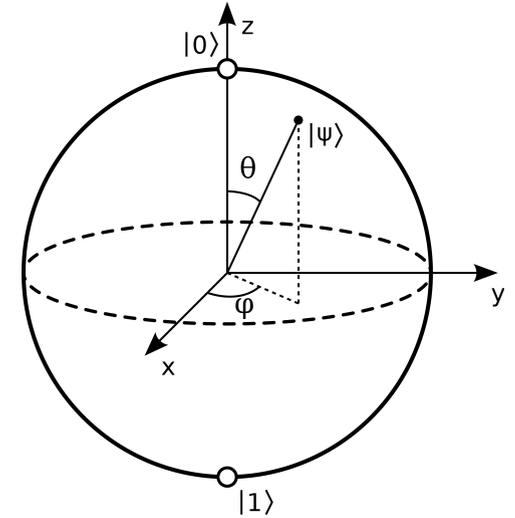
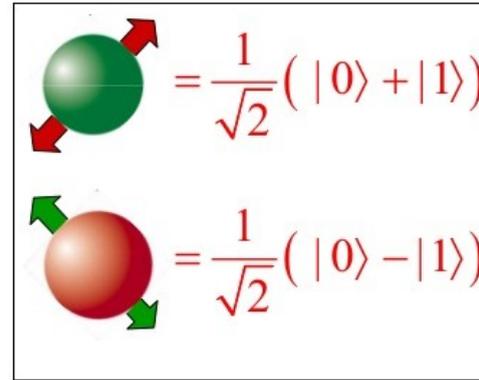
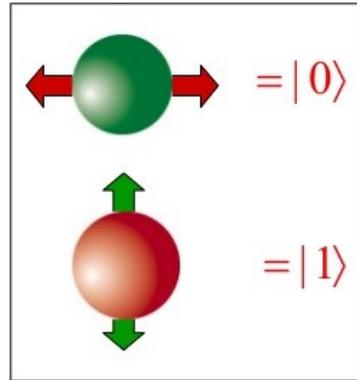


Discrete and continuous variables

DV :

information encoded in d -
level systems
(typically $d = 2$, **qubits**)

$$\alpha |0\rangle + \beta |1\rangle$$

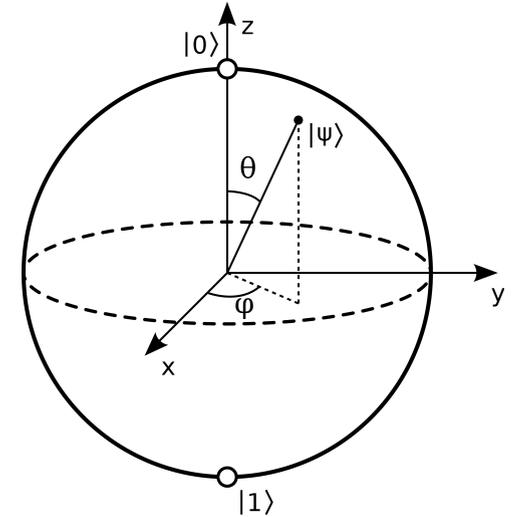
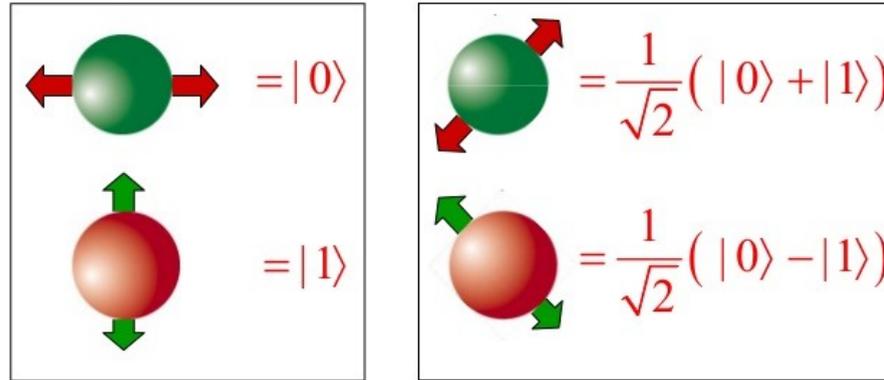


Discrete and continuous variables

DV :

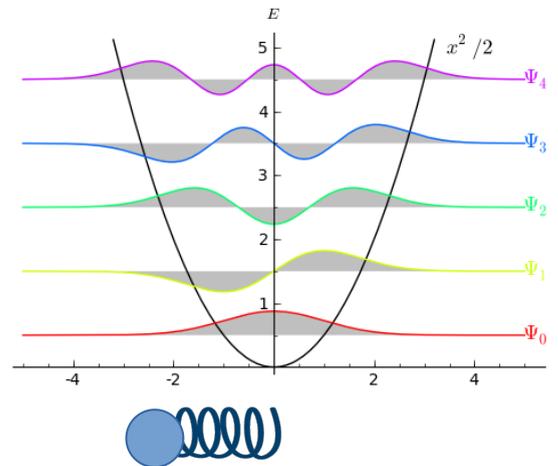
information encoded in d -level systems
(typically $d = 2$, **qubits**)

$$\alpha |0\rangle + \beta |1\rangle$$



CV :

information encoded in **oscillators**, observables with continuous spectrum,
e.g. : \hat{q} , \hat{p}

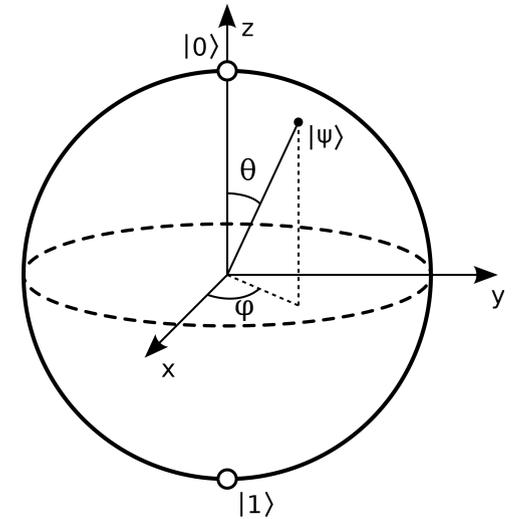
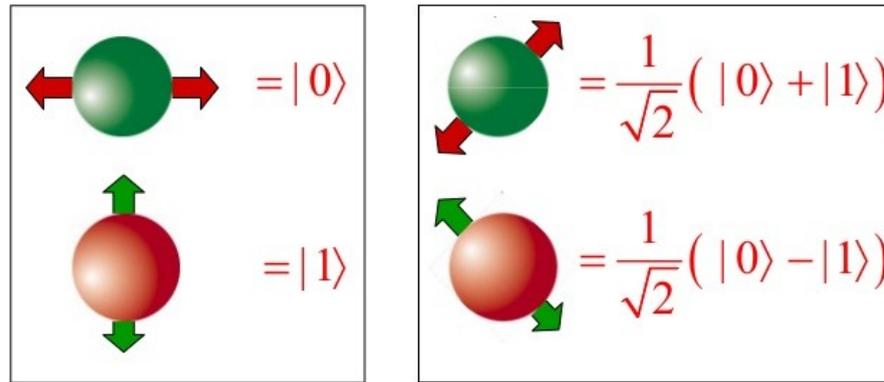


Discrete and continuous variables

DV :

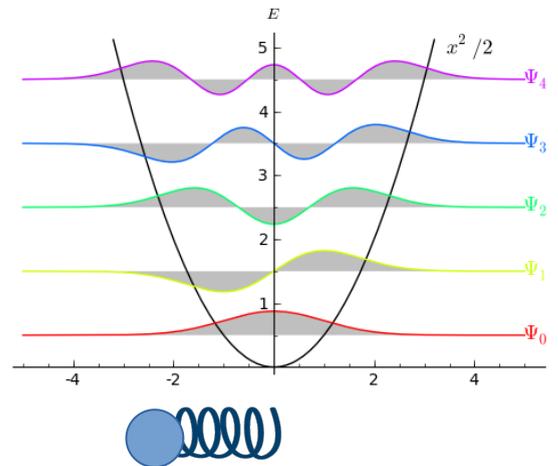
information encoded in d -level systems
(typically $d = 2$, **qubits**)

$$\alpha |0\rangle + \beta |1\rangle$$



CV :

information encoded in **oscillators**, observables with continuous spectrum,
e.g. : \hat{q} , \hat{p}



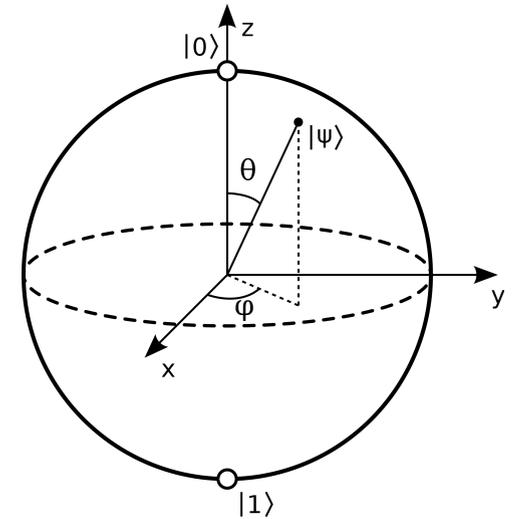
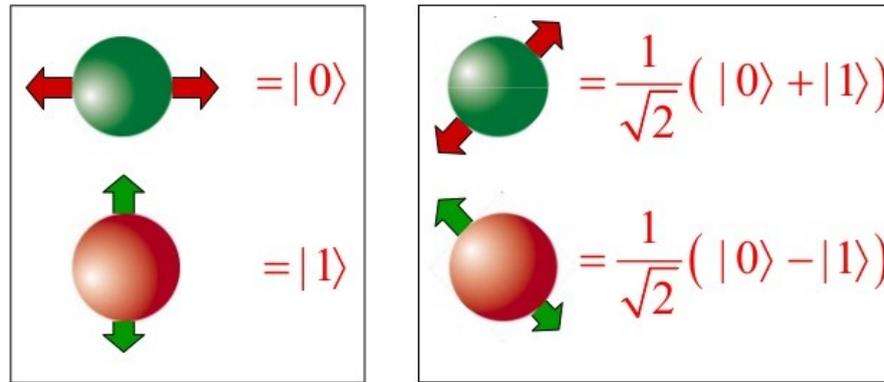
Can be light modes, LC circuits, mechanical oscillators, CoM of cold atoms...

Discrete and continuous variables

DV :

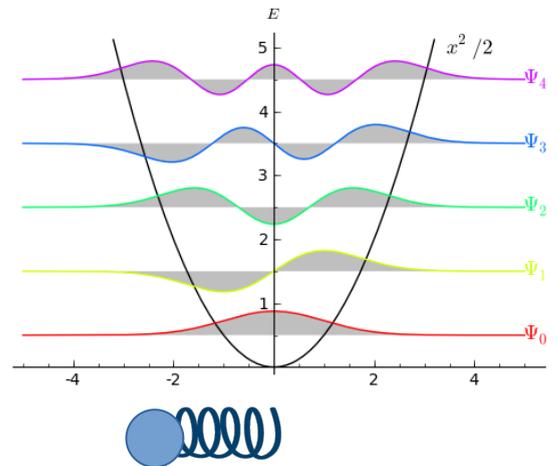
information encoded in d -level systems
(typically $d = 2$, **qubits**)

$$\alpha |0\rangle + \beta |1\rangle$$

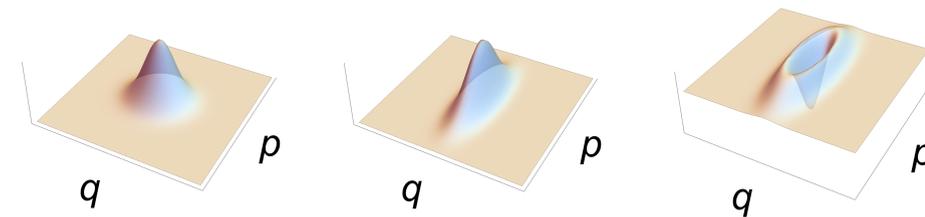


CV :

information encoded in **oscillators**, observables with continuous spectrum,
e.g. : \hat{q} , \hat{p}



In phase space:
Wigner Function



Quasi-probability distribution

Can be light modes, LC circuits, mechanical oscillators, CoM of cold atoms...

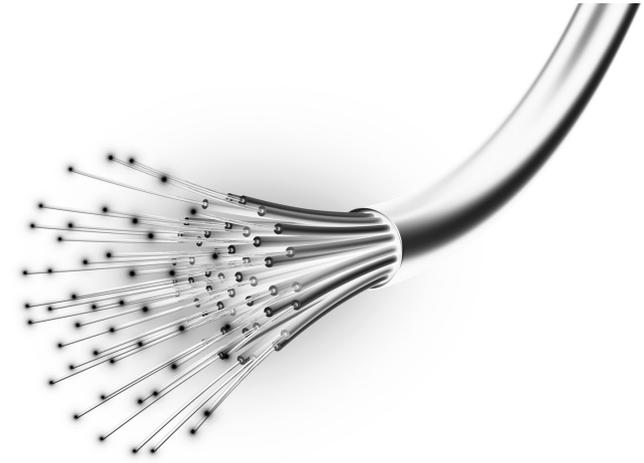
Why optics?

Why CV?

Why optics?

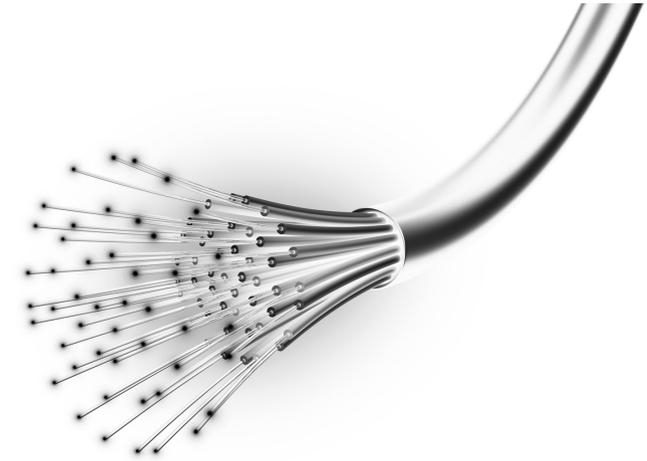
Light is used for ordinary (classical) communications

➡ Lots of know-how, technology



Why CV?

Why optics?



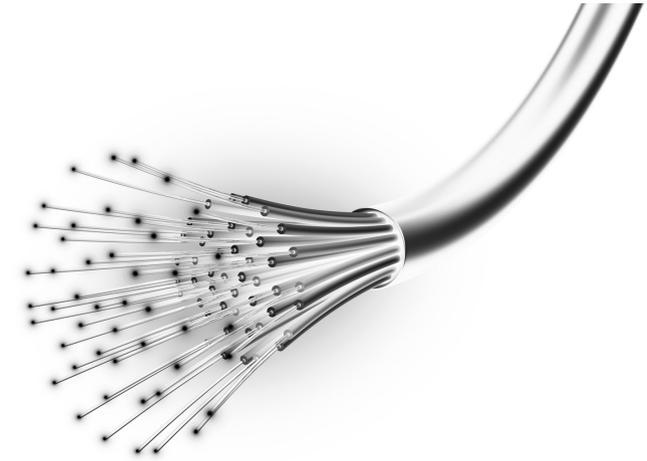
Light is used for ordinary (classical) communications

→ Lots of know-how, technology

Photons interact weakly → Easy to protect fragile quantum states

Why CV?

Why optics?



Light is used for ordinary (classical) communications

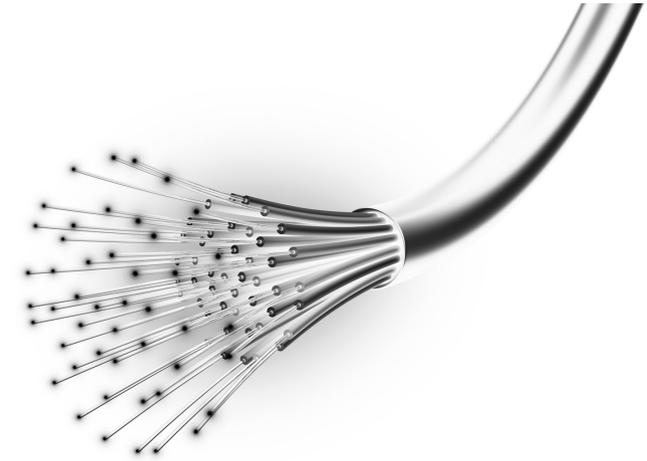
→ Lots of know-how, technology

Photons interact weakly → Easy to protect fragile quantum states

Easily scalable → If you get one, you can get plenty

Why CV?

Why optics?



Light is used for ordinary (classical) communications

➡ Lots of know-how, technology

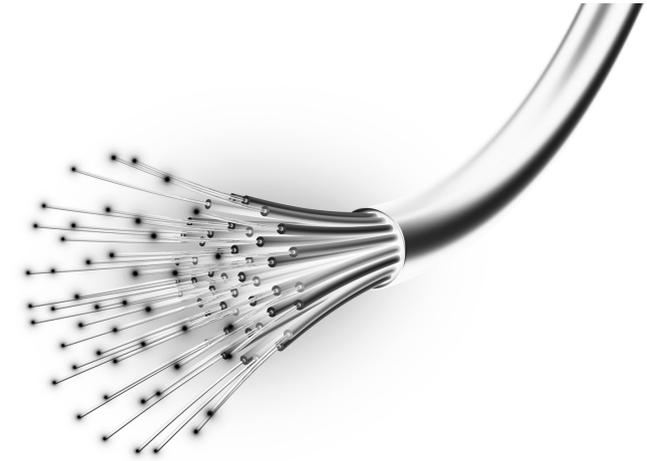
Photons interact weakly ➡ Easy to protect fragile quantum states

Easily scalable ➡ If you get one, you can get plenty

Why CV?

- Many systems *are* CV systems

Why optics?



Light is used for ordinary (classical) communications

➡ Lots of know-how, technology

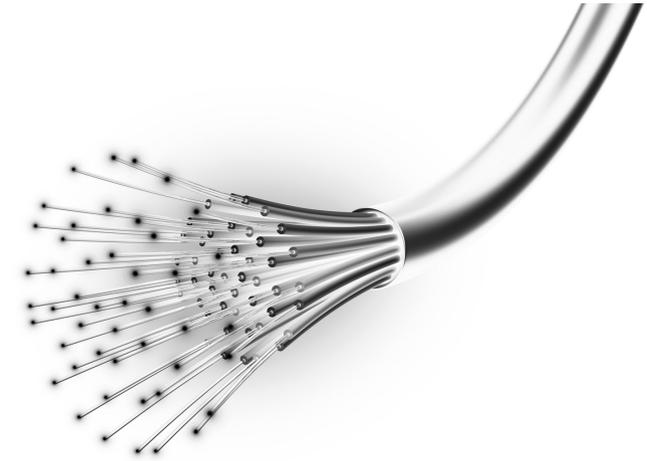
Photons interact weakly ➡ Easy to protect fragile quantum states

Easily scalable ➡ If you get one, you can get plenty

Why CV?

- Many systems *are* CV systems
- Deterministic, massively multimode entangled systems (with much flexibility)

Why optics?



Light is used for ordinary (classical) communications

➡ Lots of know-how, technology

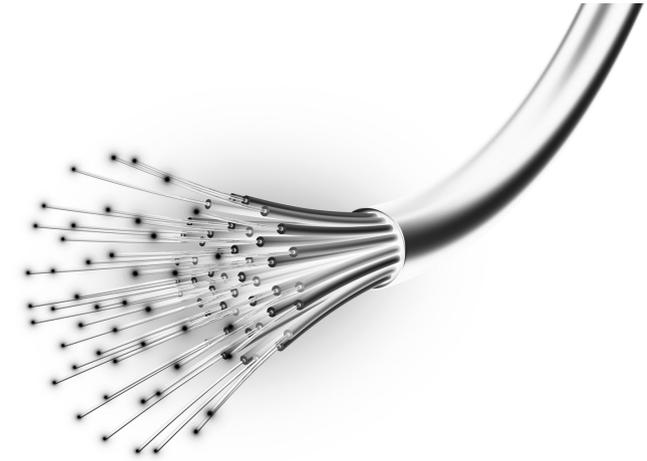
Photons interact weakly ➡ Easy to protect fragile quantum states

Easily scalable ➡ If you get one, you can get plenty

Why CV?

- Many systems *are* CV systems
- Deterministic, massively multimode entangled systems (with much flexibility)
- Loss-resistant (ECC codes)

Why optics?



Light is used for ordinary (classical) communications

➡ Lots of know-how, technology

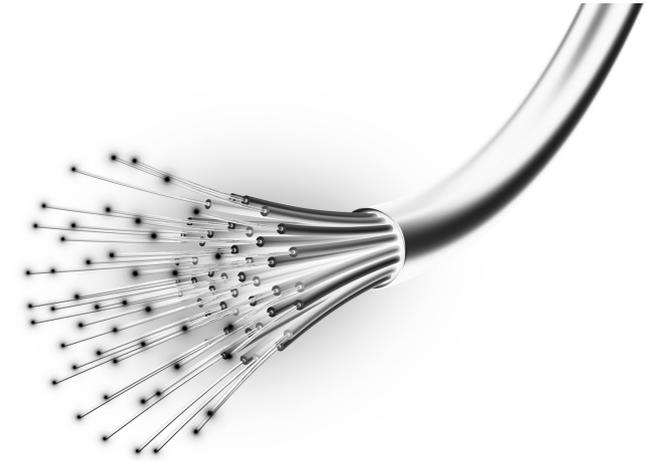
Photons interact weakly ➡ Easy to protect fragile quantum states

Easily scalable ➡ If you get one, you can get plenty

Why CV?

- Many systems *are* CV systems
- Deterministic, massively multimode entangled systems (with much flexibility)
- Loss-resistant (ECC codes)
- Complementary “easy” operations with respect to other systems (hybrid devices)

Why optics?



Light is used for ordinary (classical) communications

➡ Lots of know-how, technology

Photons interact weakly ➡ Easy to protect fragile quantum states

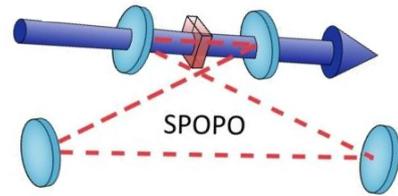
Easily scalable ➡ If you get one, you can get plenty

Why CV?

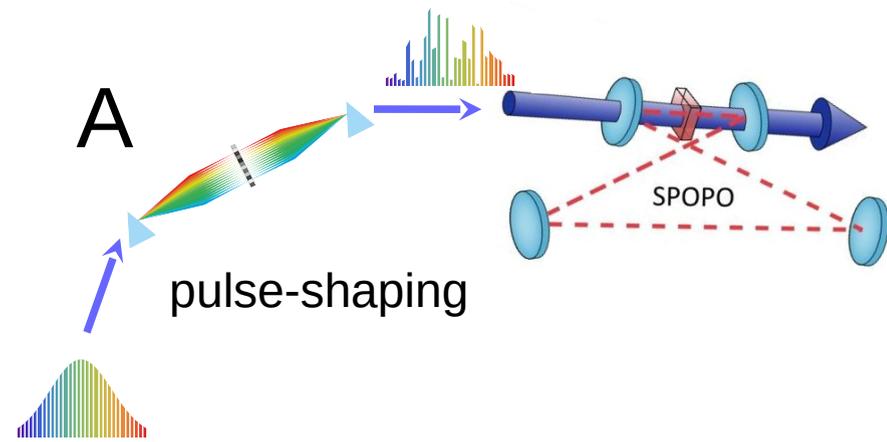
- Many systems *are* CV systems
- Deterministic, massively multimode entangled systems (with much flexibility)
- Loss-resistant (ECC codes)
- Complementary “easy” operations with respect to other systems (hybrid devices)
- New sets of problems (Boson sampling)

Past research

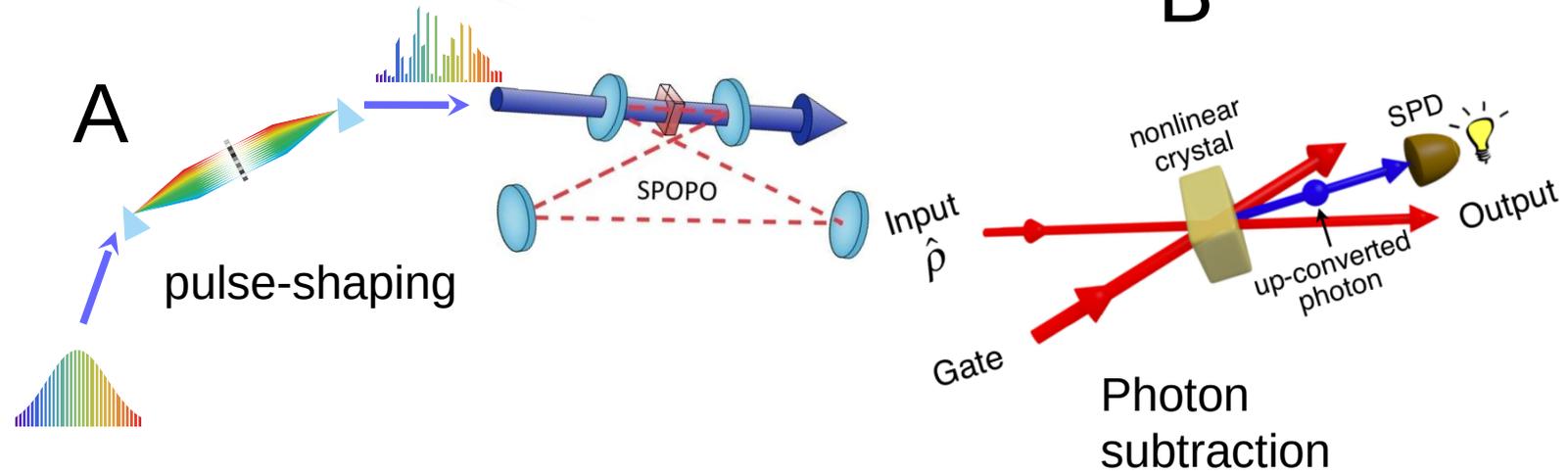
Overview



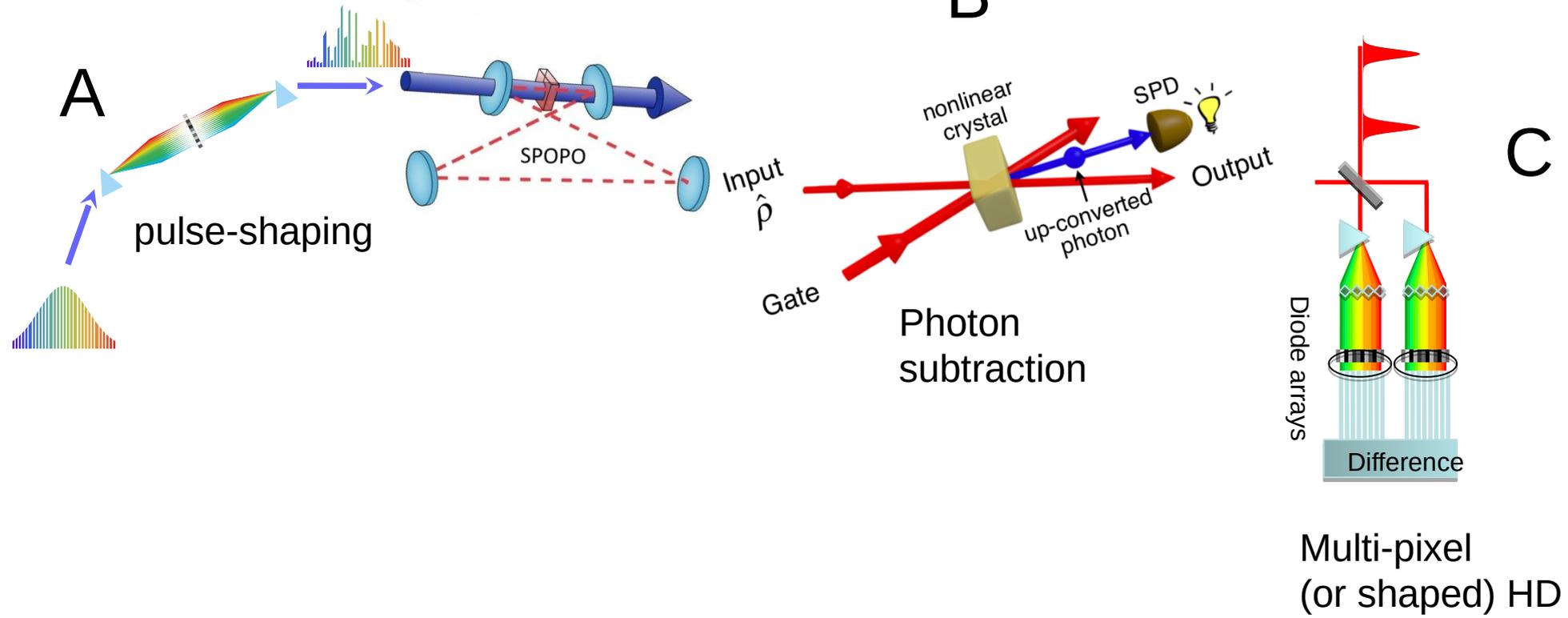
Overview



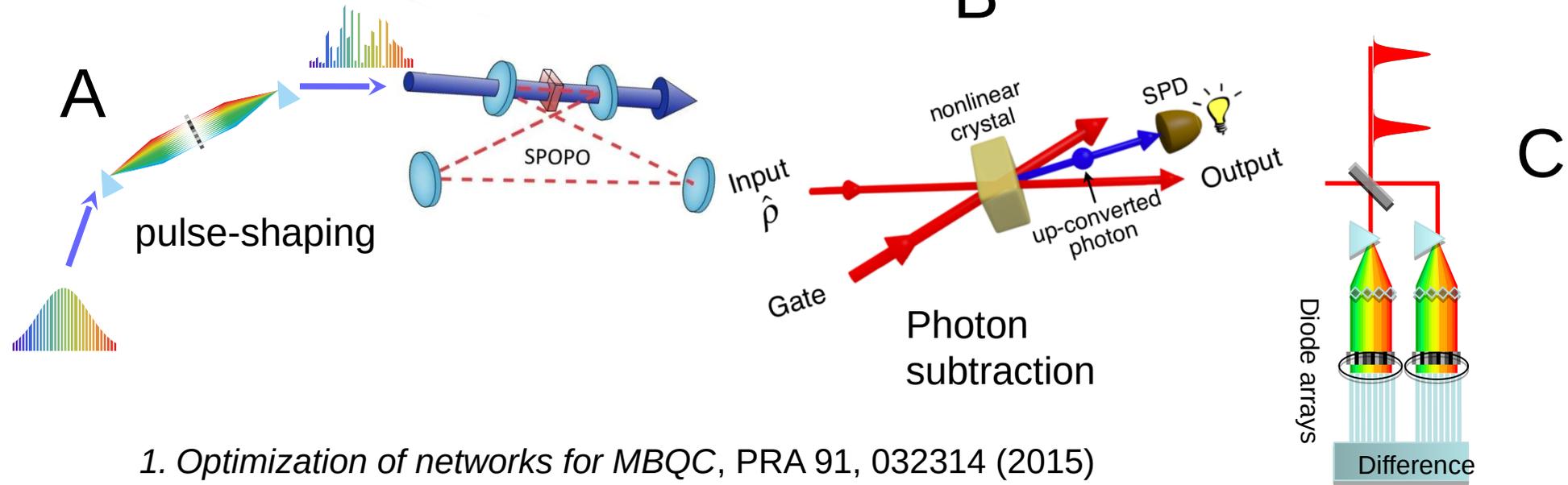
Overview



Overview



Overview



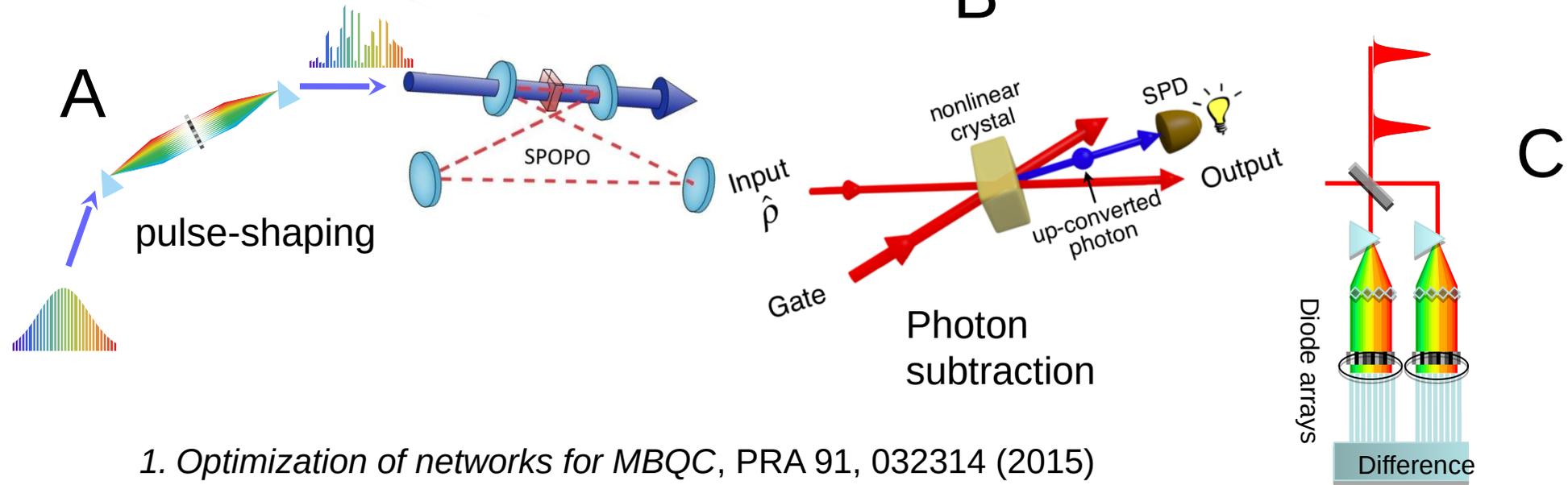
1. *Optimization of networks for MBQC*, PRA 91, 032314 (2015)
2. *Direct approach to Gaussian MBQC*, PRA 94, 062332 (2016)
3. *Multimode entanglement in reconfigurable graph states using optical frequency combs*, NatComm 8, 15645 (2017)
4. *Polynomial approximation of non-Gaussian unitaries by counting one photon at a time*, PRA 95 (5), 052352 (2017)
5. *Violating Bell inequalities with entangled optical frequency combs and multipixel homodyne detection*, PRA 98, 062101 (2018)
6. *High-dimensional quantum encoding via photon-subtracted squeezed states*, PRA A 99, 022342 (2019)
7. *Versatile engineering of multimode squeezed states by optimizing the pump spectral profile in spontaneous parametric down-conversion*, PRA 97, 033808 (2018)
8. *Reconfigurable optical implementation of quantum complex networks*, NJP 20, 053024 (2018)
9. *Bloch-Messiah reduction for twin beams of light*, PRA 100, 013837 (2019)
10. *Random coding for sharing bosonic quantum secrets*, PRA 100, 022303 (2019)

C

B (+C)

A+C

Overview



1. *Optimization of networks for MBQC*, PRA 91, 032314 (2015)
2. *Direct approach to Gaussian MBQC*, PRA 94, 062332 (2016)
3. ***Multimode entanglement in reconfigurable graph states using optical frequency combs***, NatComm 8, 15645 (2017)
4. ***Polynomial approximation of non-Gaussian unitaries by counting one photon at a time***, PRA 95 (5), 052352 (2017)
5. *Violating Bell inequalities with entangled optical frequency combs and multipixel homodyne detection*, PRA 98, 062101 (2018)
6. *High-dimensional quantum encoding via photon-subtracted squeezed states*, PRA A 99, 022342 (2019)
7. ***Versatile engineering of multimode squeezed states by optimizing the pump spectral profile in spontaneous parametric down-conversion***, PRA 97, 033808 (2018)
8. ***Reconfigurable optical implementation of quantum complex networks***, NJP 20, 053024 (2018)
9. *Bloch-Messiah reduction for twin beams of light*, PRA 100, 013837 (2019)
10. ***Random coding for sharing bosonic quantum secrets***, PRA 100, 022303 (2019)

C

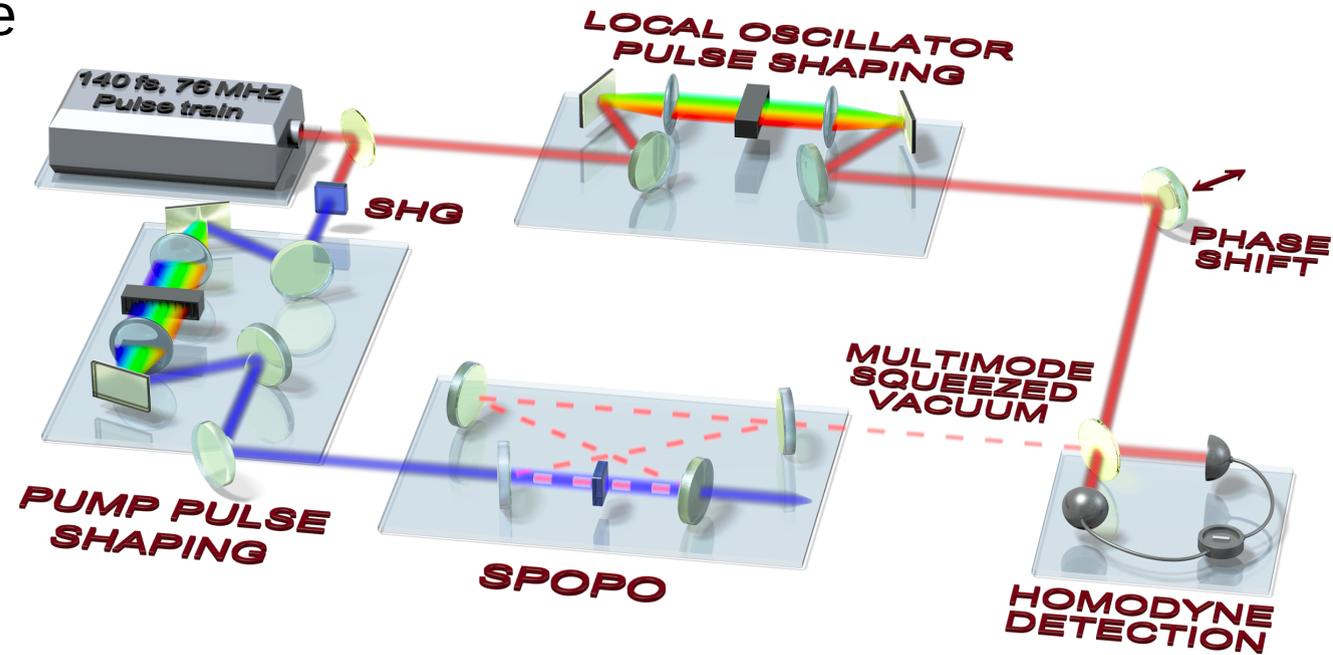
B (+C)

A+C

Entanglement engineering

Image by
Valérian Thiel

Motivation: Experimental setup can generate multi-mode, entangled states of light. Such states can be resources for quantum information processing.

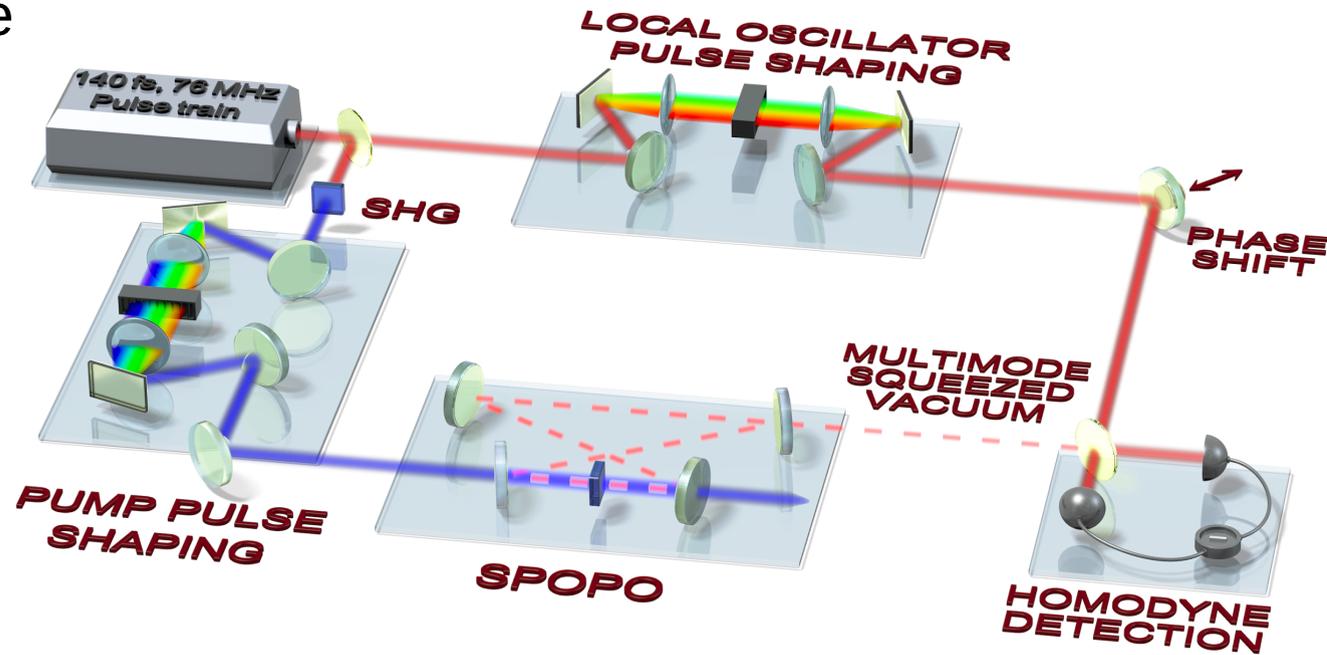


Entanglement engineering

Image by
Valérian Thiel

Motivation: Experimental setup can generate multi-mode, entangled states of light. Such states can be resources for quantum information processing.

Question: is the setup adapted to QIP and if not, how to change it?

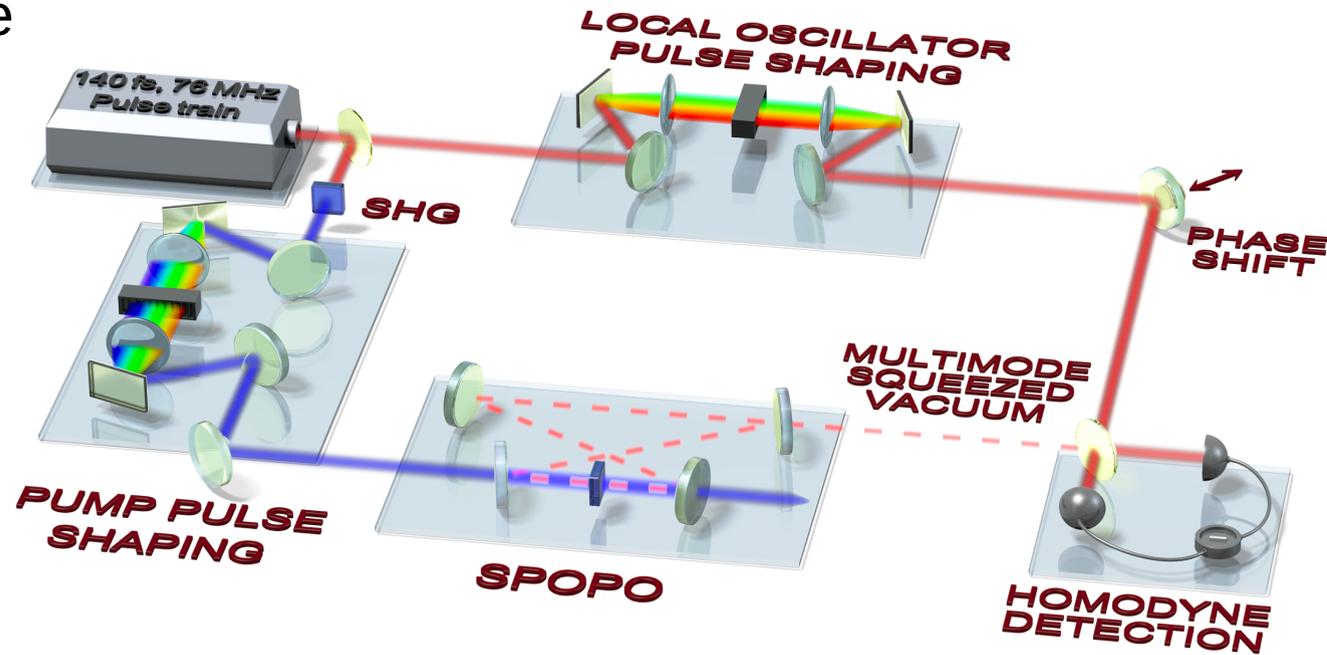
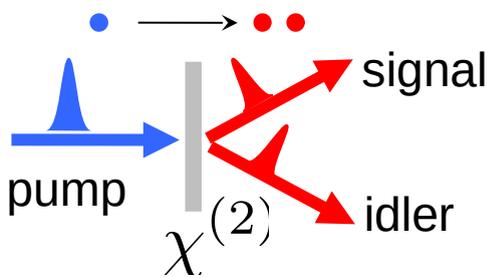


Entanglement engineering

Image by
Valérian Thiel

Motivation: Experimental setup can generate multi-mode, entangled states of light. Such states can be resources for quantum information processing.

Question: is the setup adapted to QIP and if not, how to change it?

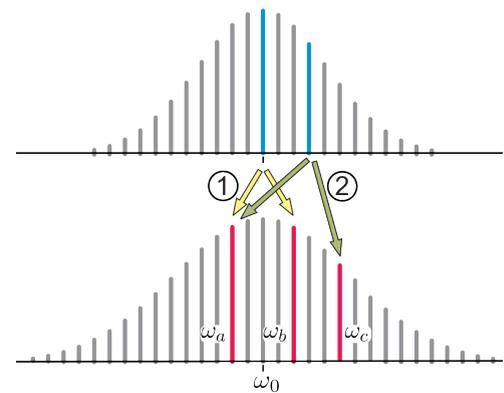
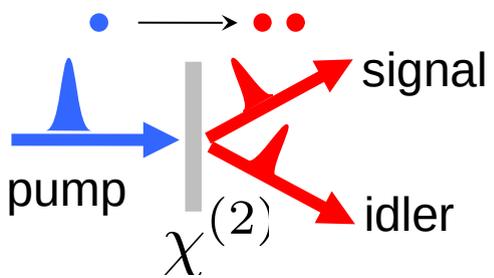


Entanglement engineering

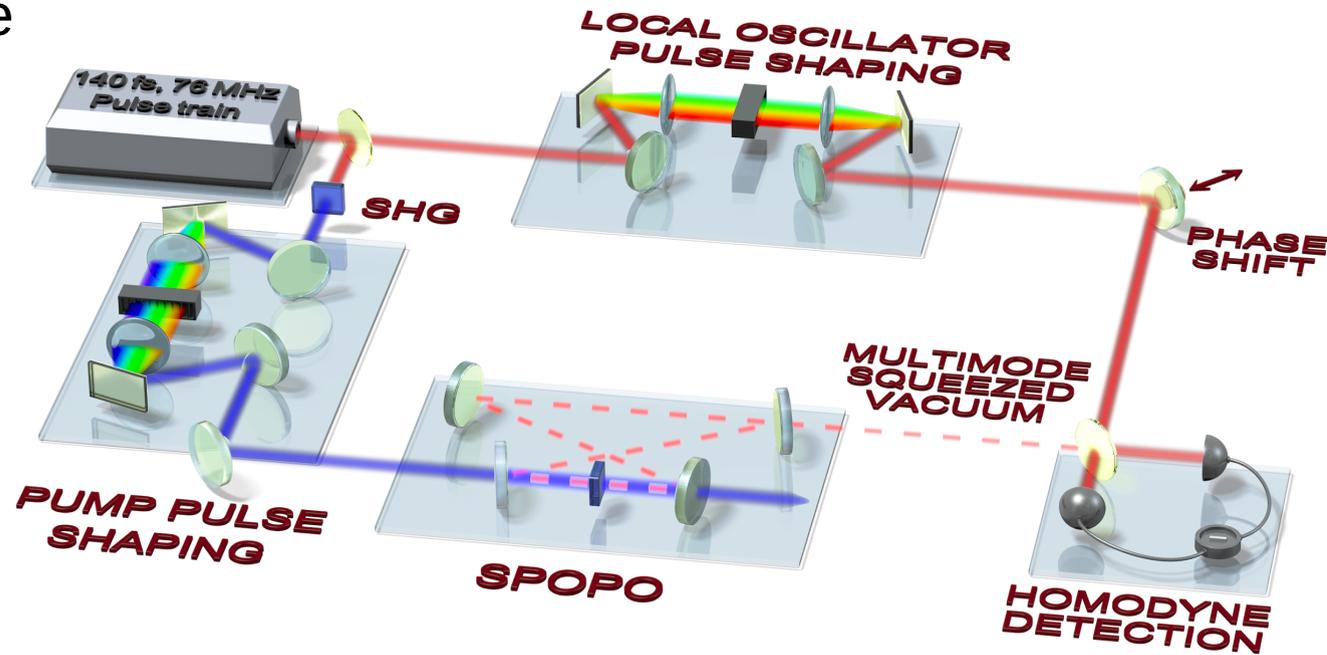
Image by Valérian Thiel

Motivation: Experimental setup can generate multi-mode, entangled states of light. Such states can be resources for quantum information processing.

Question: is the setup adapted to QIP and if not, how to change it?



Entanglement from spontaneous parametric down-conversion of optical frequency combs

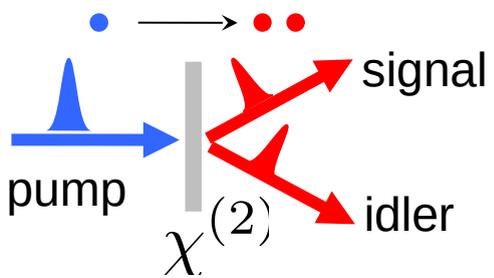
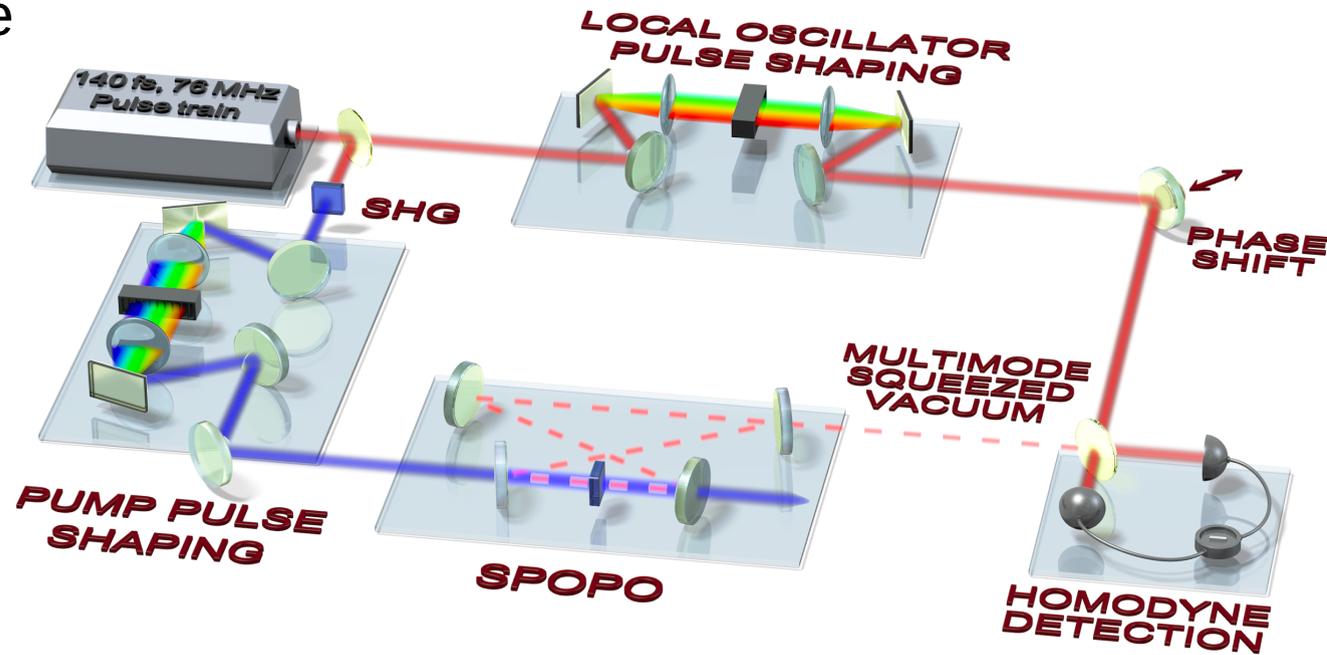


Entanglement engineering

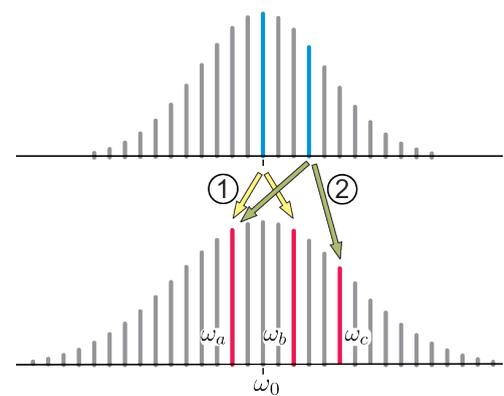
Image by Valérian Thiel

Motivation: Experimental setup can generate multi-mode, entangled states of light. Such states can be resources for quantum information processing.

Question: is the setup adapted to QIP and if not, how to change it?



Entanglement from spontaneous parametric down-conversion of optical frequency combs



Effective quadratic Hamiltonian:

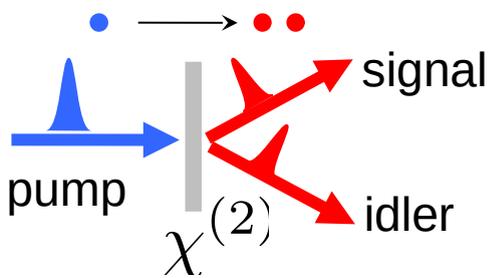
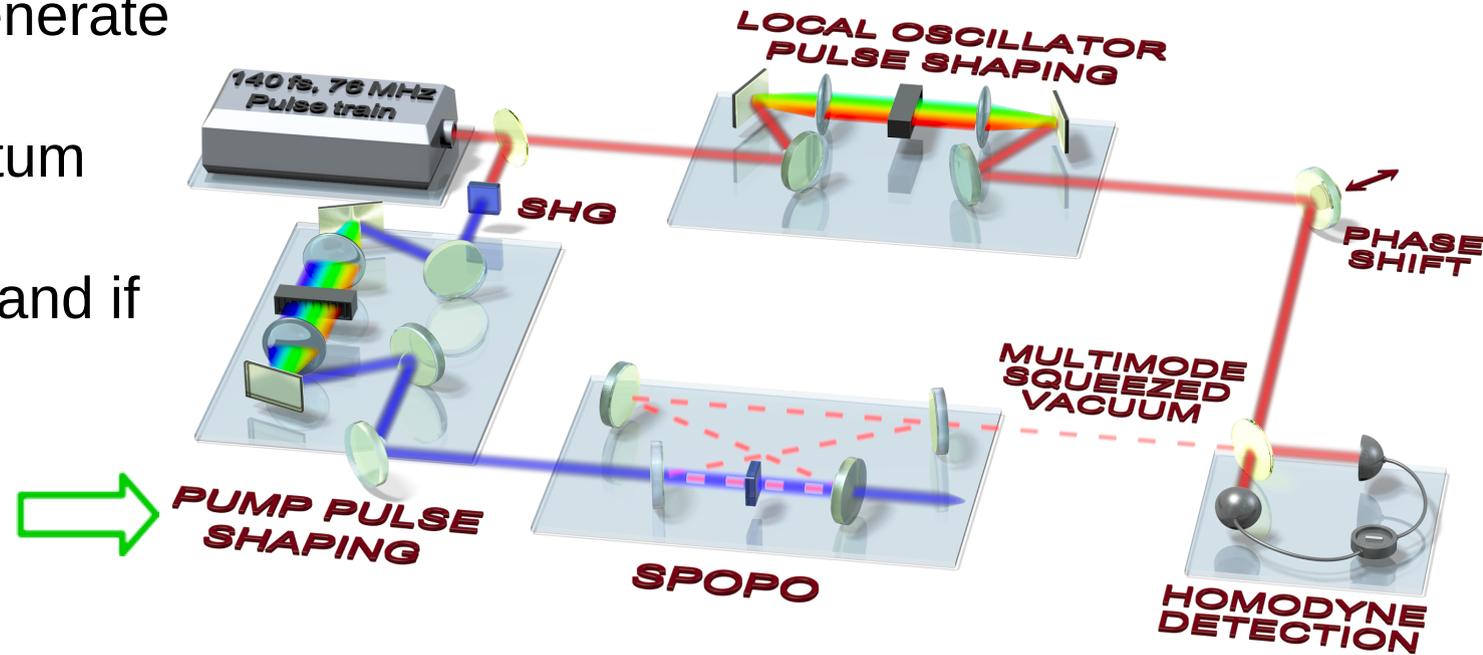
$$H = i \sum_{m,q} \mathcal{L}_{m,q} \hat{a}_{\omega_m}^\dagger \hat{a}_{\omega_q}^\dagger + \text{h.c.}$$

Entanglement engineering

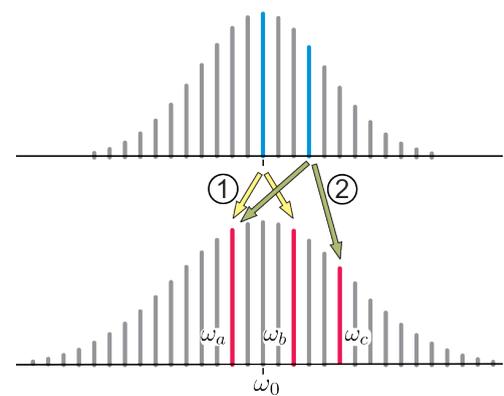
Image by
Valérian Thiel

Motivation: Experimental setup can generate multi-mode, entangled states of light. Such states can be resources for quantum information processing.

Question: is the setup adapted to QIP and if not, how to change it?



Entanglement from spontaneous parametric down-conversion of optical frequency combs



Effective quadratic Hamiltonian:

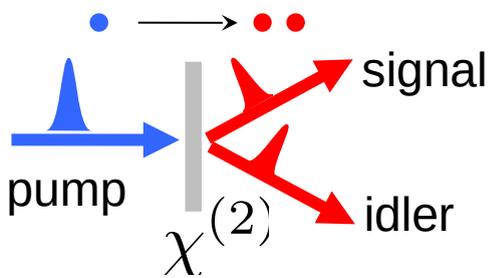
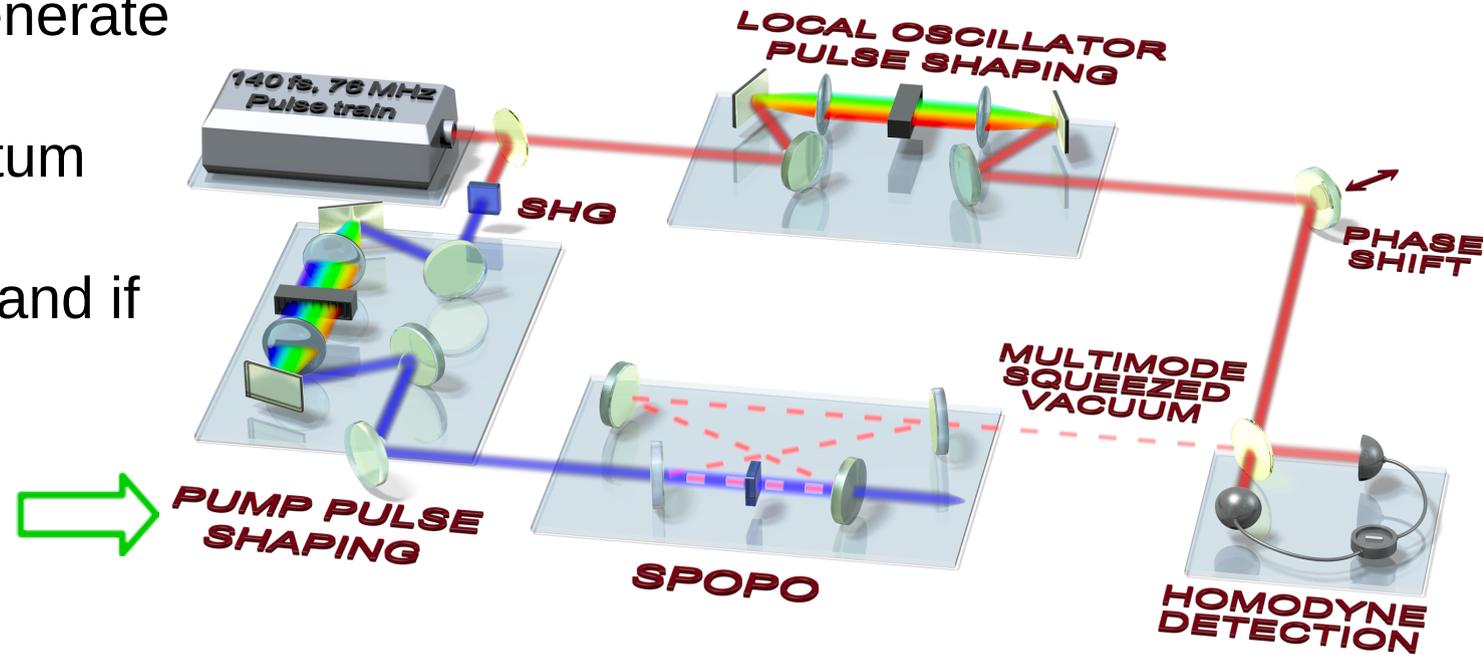
$$H = i \sum_{m,q} \mathcal{L}_{m,q} \hat{a}_{\omega_m}^\dagger \hat{a}_{\omega_q}^\dagger + \text{h.c.}$$

Entanglement engineering

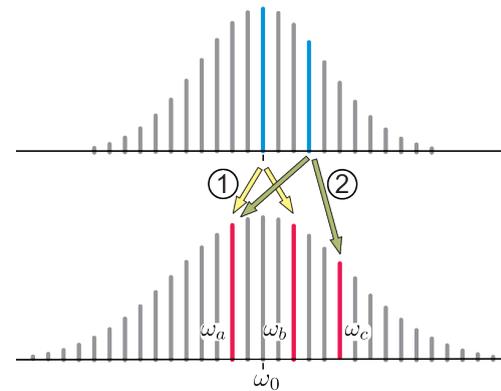
Image by Valérian Thiel

Motivation: Experimental setup can generate multi-mode, entangled states of light. Such states can be resources for quantum information processing.

Question: is the setup adapted to QIP and if not, how to change it?



Entanglement from spontaneous parametric down-conversion of optical frequency combs



Effective quadratic Hamiltonian:

$$H = i \sum_{m,q} \mathcal{L}_{m,q} \hat{a}_{\omega_m}^\dagger \hat{a}_{\omega_q}^\dagger + \text{h.c.}$$

Our contribution:

Combining pump-shape and detection system optimization it is possible to realize QIP including

- Measurement-based Q comp.
- Simulation of complex networks (Turku)
- Secret sharing (more later)

Non-Gaussianity from single-photon detection

Motivation:

- 1.No Q Advantage without non-Gaussian
- 2.Realizable non-Gauss: single photon ops

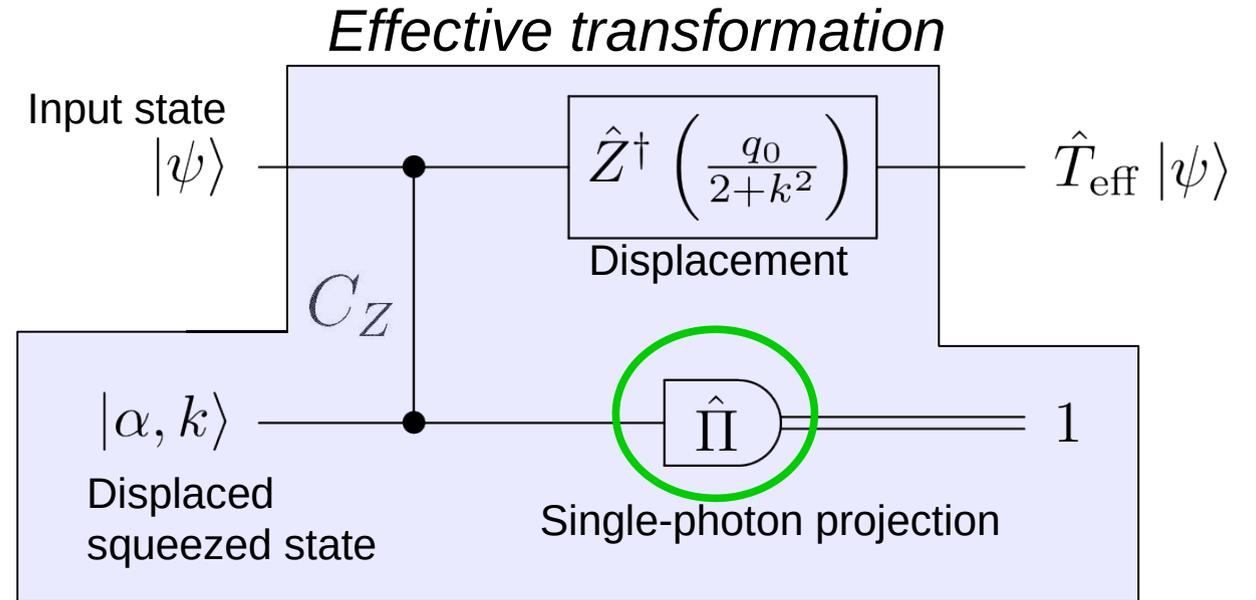
Non-Gaussianity from single-photon detection

Motivation:

- 1.No Q Advantage without non-Gaussian
- 2.Realizable non-Gauss: single photon ops

Strategy:

- 1.Entangle input to a Gaussian state
- 2.Detect a single photon (probabilistic)
- 3.Perform correction
- 4.Repeat



$$\hat{T}_{\text{eff}} = \tilde{\mathcal{N}} \exp \left\{ - \left(\frac{k^2}{4 + 2k^2} \right) (\hat{q} + p_0)^2 \right\} \left(\hat{q} - \lambda(\alpha, k) \right)$$

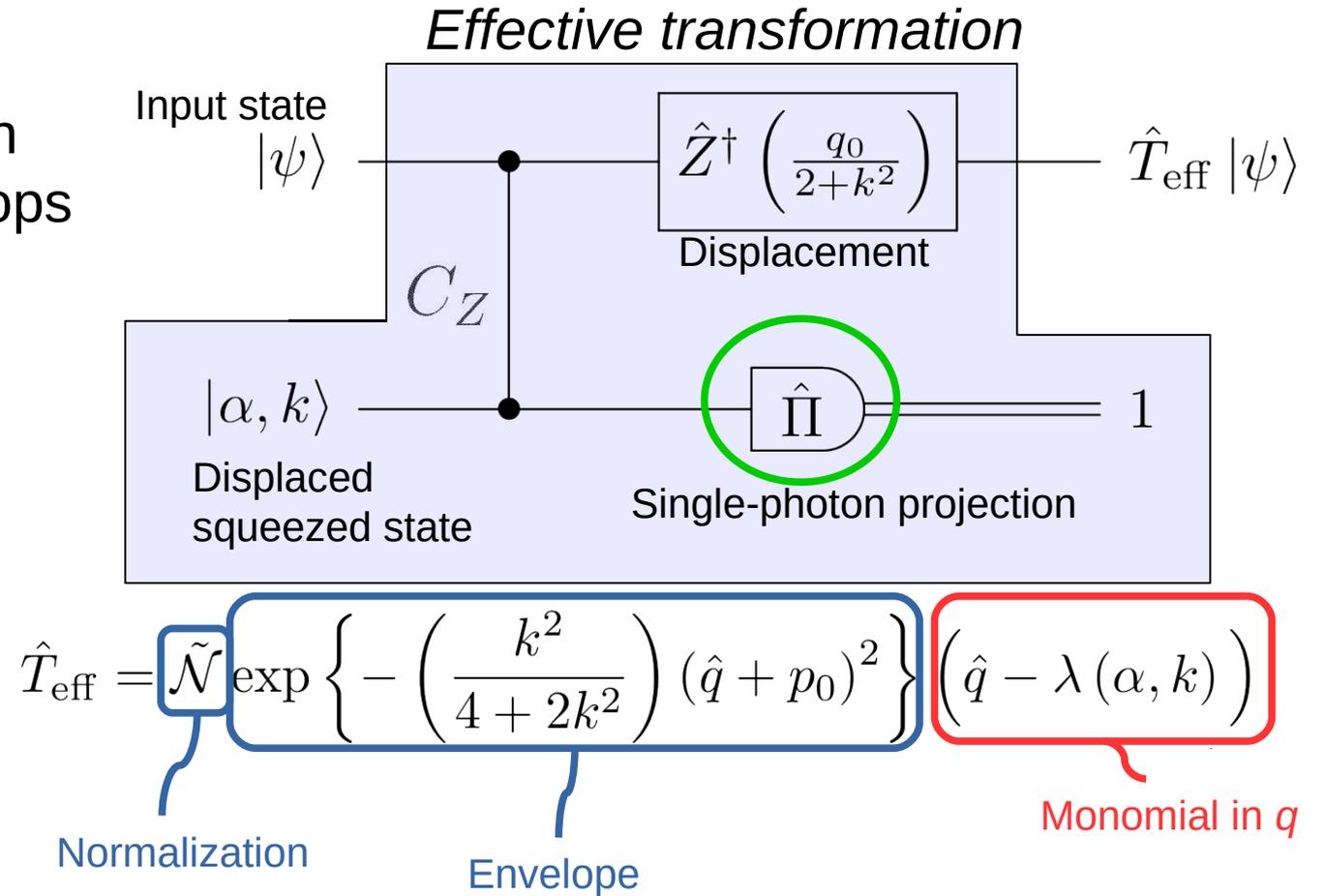
Non-Gaussianity from single-photon detection

Motivation:

- 1.No Q Advantage without non-Gaussian
- 2.Realizable non-Gauss: single photon ops

Strategy:

- 1.Entangle input to a Gaussian state
- 2.Detect a single photon (probabilistic)
- 3.Perform correction
- 4.Repeat



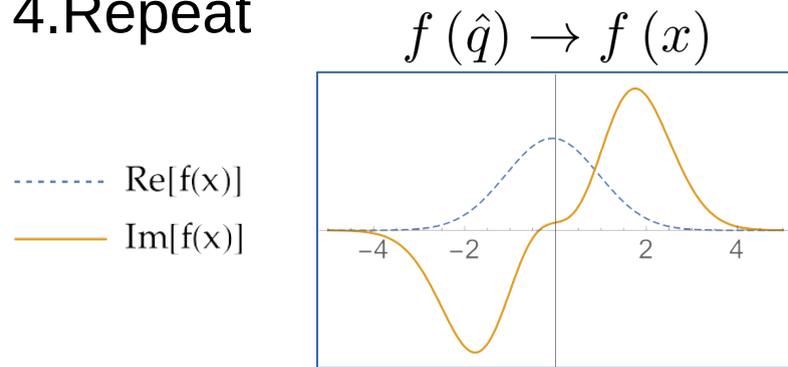
Non-Gaussianity from single-photon detection

Motivation:

- 1.No Q Advantage without non-Gaussian
- 2.Realizable non-Gauss: single photon ops

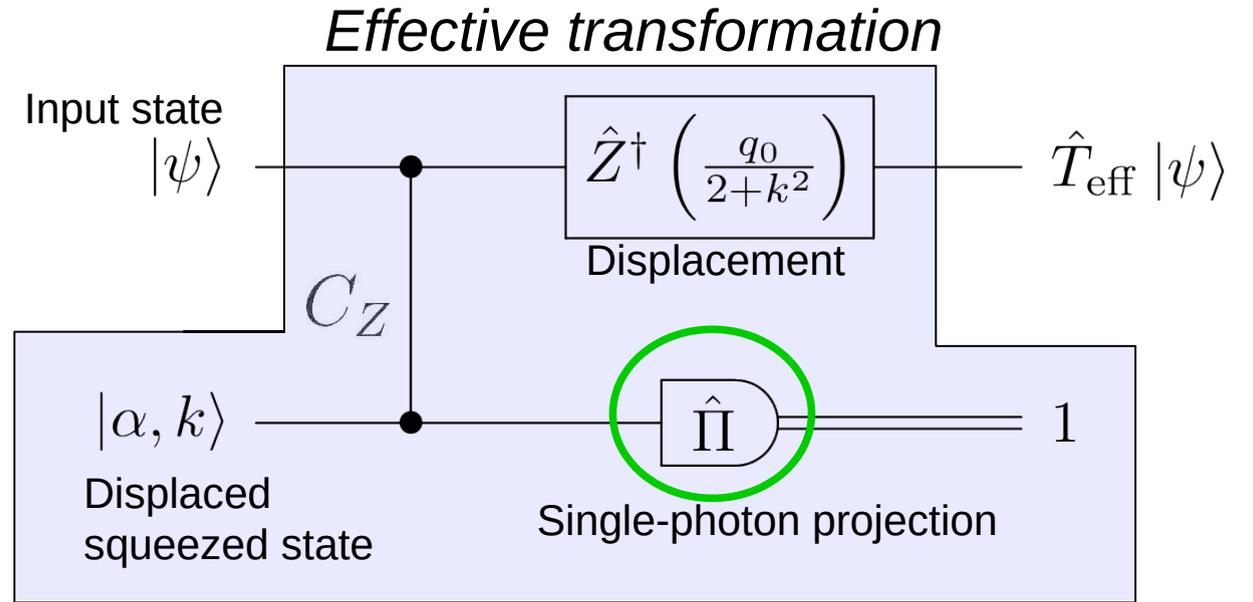
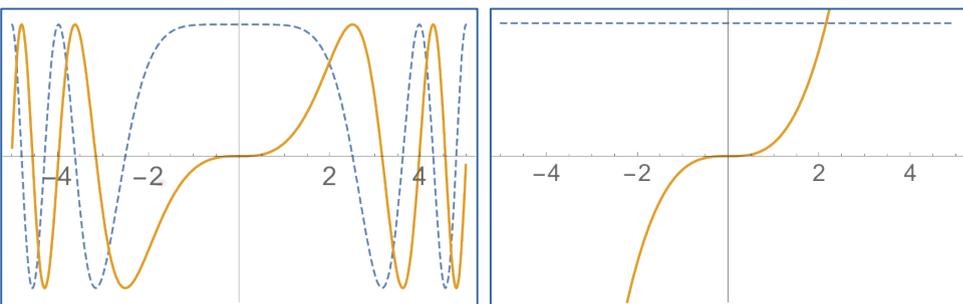
Strategy:

- 1.Entangle input to a Gaussian state
- 2.Detect a single photon (probabilistic)
- 3.Perform correction
- 4.Repeat



$$e^{0.1ix^3}$$

$$1 + 0.1ix^3$$



$$\hat{T}_{\text{eff}} = \underbrace{\tilde{\mathcal{N}}}_{\text{Normalization}} \exp\left\{-\underbrace{\left(\frac{k^2}{4+2k^2}\right)}_{\text{Envelope}} (\hat{q} + p_0)^2\right\} \underbrace{\left(\hat{q} - \lambda(\alpha, k)\right)}_{\text{Monomial in } q}$$

$$e^{i\nu\hat{q}^3} \approx \mathbb{I} + i\nu\hat{q}^3 = (\hat{q} - \lambda_1)(\hat{q} - \lambda_2)(\hat{q} - \lambda_3)$$

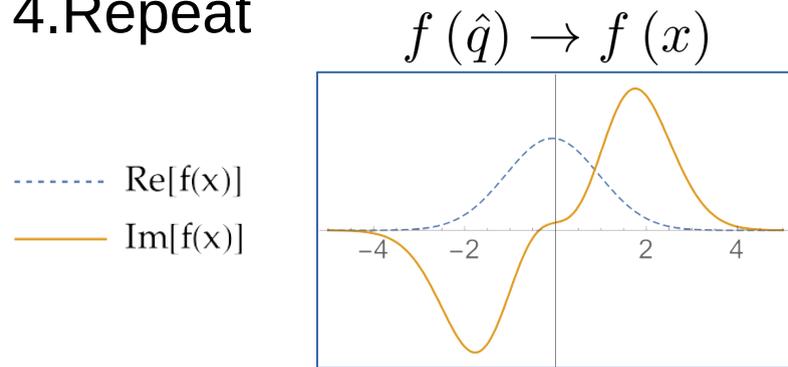
Non-Gaussianity from single-photon detection

Motivation:

- 1.No Q Advantage without non-Gaussian
- 2.Realizable non-Gauss: single photon ops

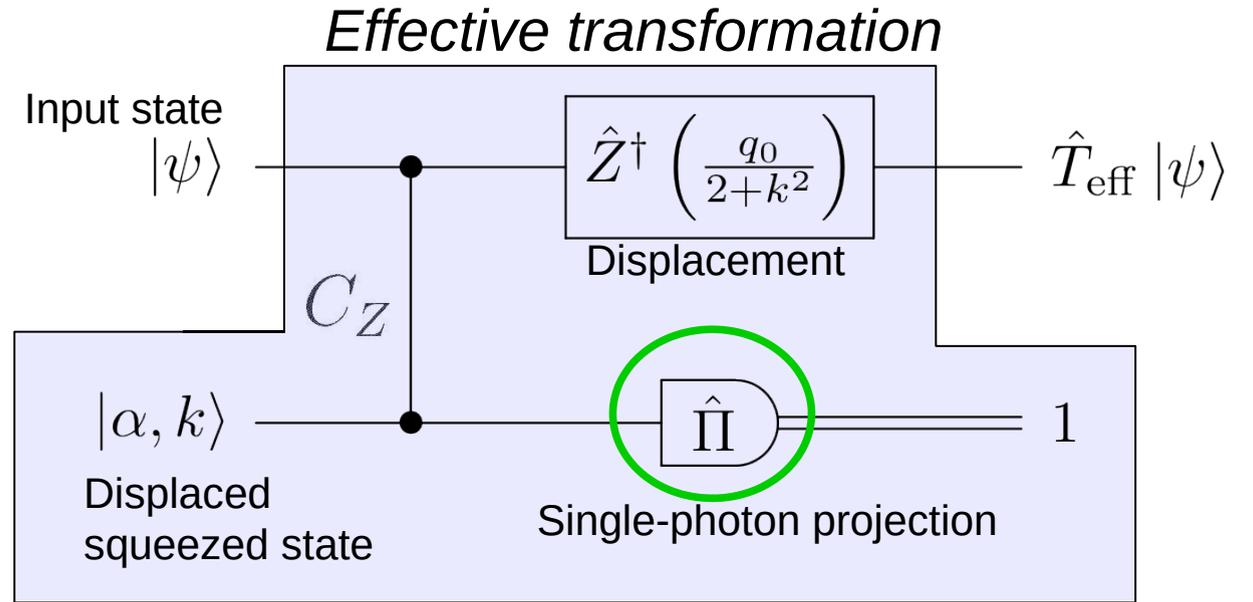
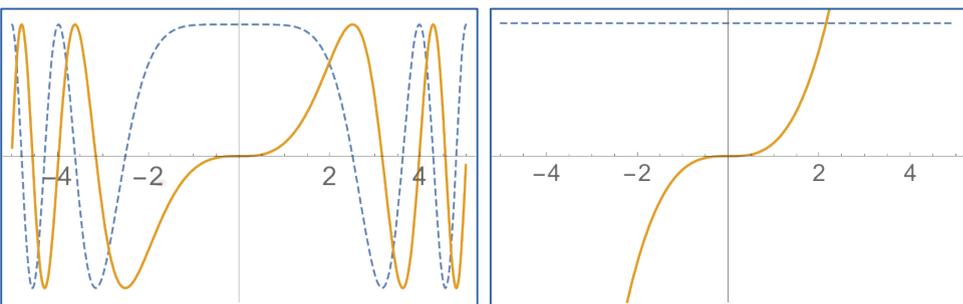
Strategy:

- 1.Entangle input to a Gaussian state
- 2.Detect a single photon (probabilistic)
- 3.Perform correction
- 4.Repeat



$$e^{0.1ix^3}$$

$$1 + 0.1ix^3$$



$$\hat{T}_{\text{eff}} = \underbrace{\tilde{\mathcal{N}}}_{\text{Normalization}} \exp \left\{ - \left(\frac{k^2}{4 + 2k^2} \right) (\hat{q} + p_0)^2 \right\} \underbrace{\left(\hat{q} - \lambda(\alpha, k) \right)}_{\text{Monomial in } q}$$

$$e^{i\nu\hat{q}^3} \approx \mathbb{I} + i\nu\hat{q}^3 = (\hat{q} - \lambda_1)(\hat{q} - \lambda_2)(\hat{q} - \lambda_3)$$

Our contribution:

Single-photon non-unitary operations can be used to approximate non-Gaussian unitary evolution

CV quantum state sharing: random codes

A **dealer** shares a **secret** with several **players** such that only **authorized subsets** of players can retrieve it if they **collaborate**

QQ: The secret is a quantum state



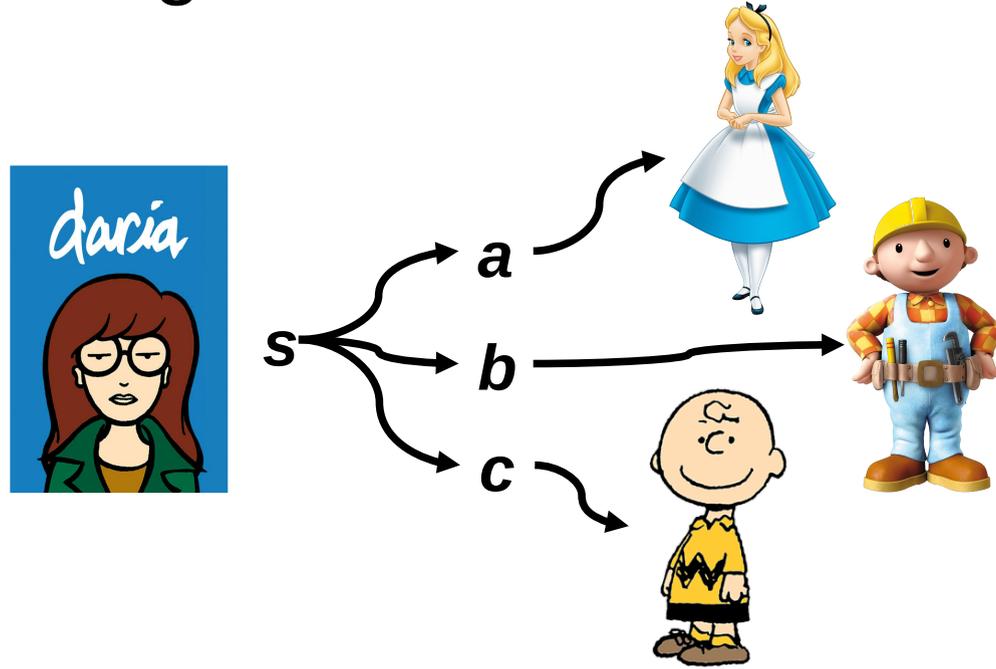
S



CV quantum state sharing: random codes

A **dealer** shares a **secret** with several **players** such that only **authorized subsets** of players can retrieve it if they **collaborate**

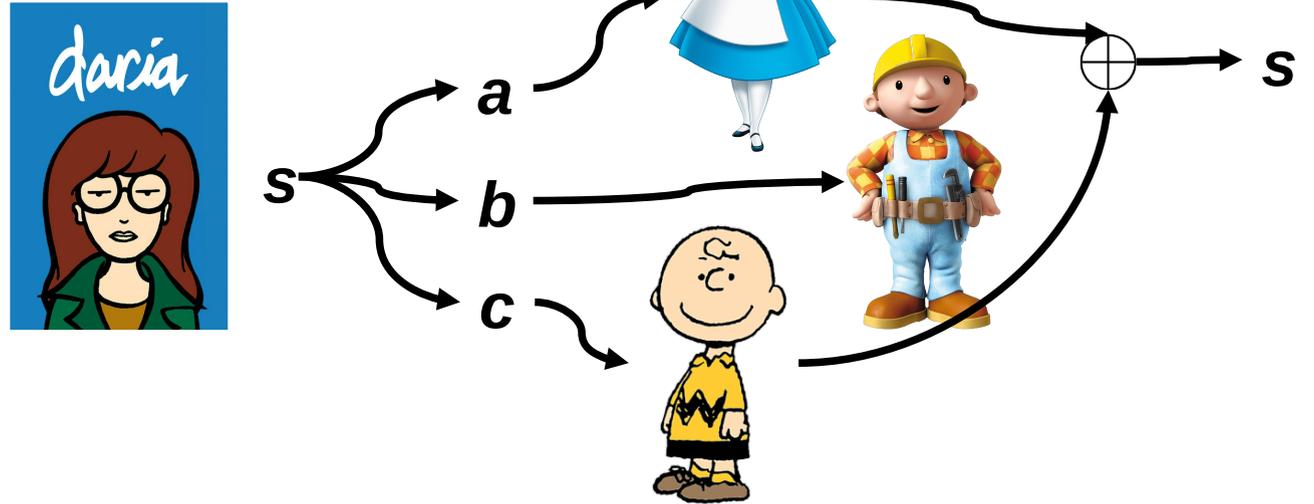
QQ: The secret is a quantum state



CV quantum state sharing: random codes

A **dealer** shares a **secret** with several **players** such that only **authorized subsets** of players can retrieve it if they **collaborate**

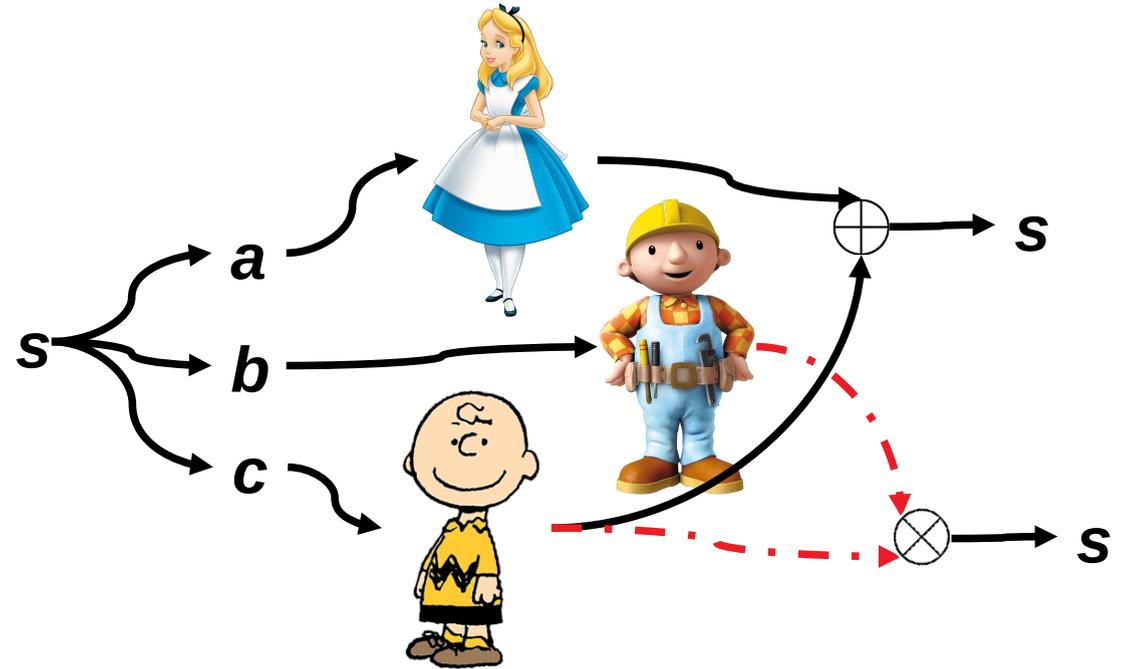
QQ: The secret is a quantum state



CV quantum state sharing: random codes

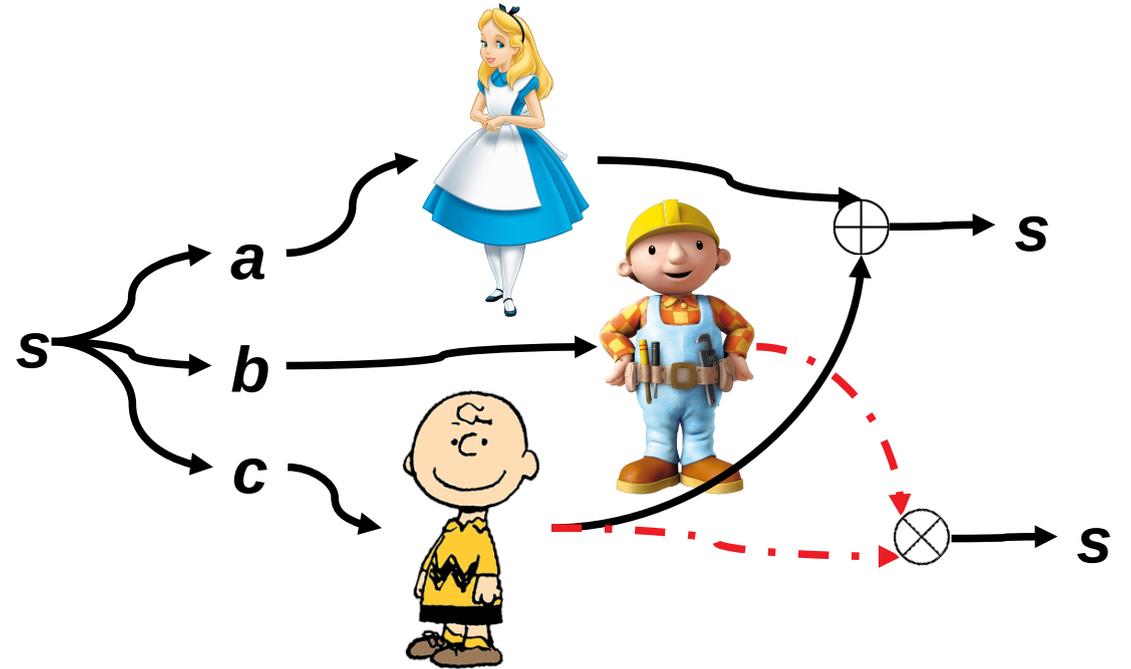
A **dealer** shares a **secret** with several **players** such that only **authorized subsets** of players can retrieve it if they **collaborate**

QQ: The secret is a quantum state



CV quantum state sharing: random codes

A **dealer** shares a **secret** with several **players** such that only **authorized subsets** of players can retrieve it if they **collaborate**

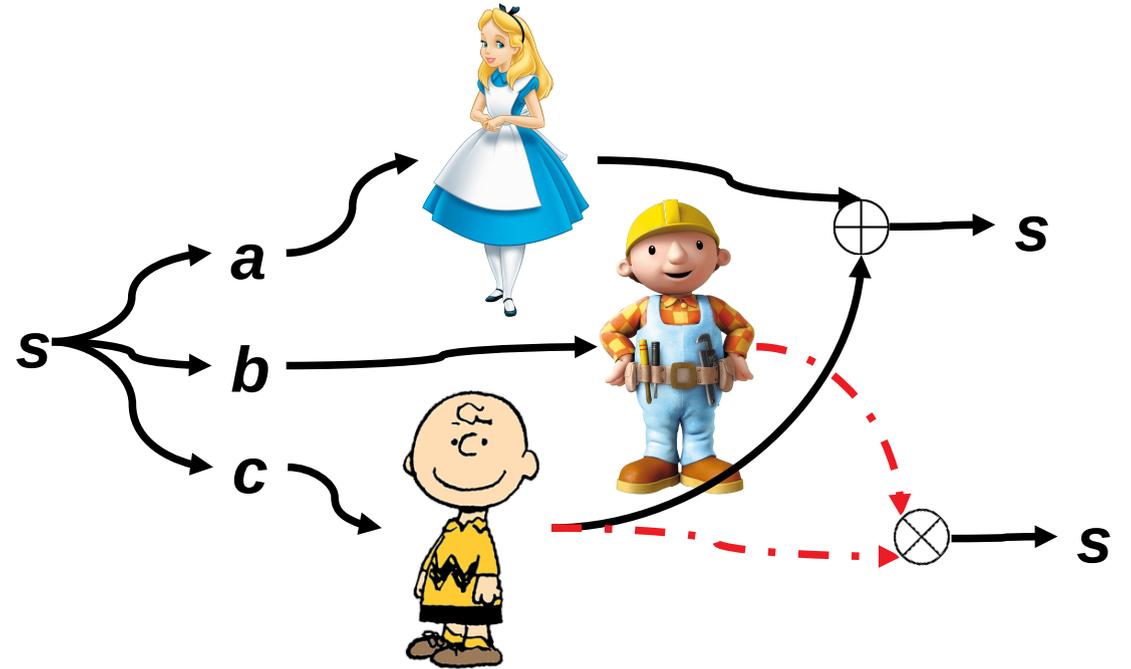


QQ: The secret is a quantum state

How it started: NatComm 8, 15645 (2017)

CV quantum state sharing: random codes

A **dealer** shares a **secret** with several **players** such that only **authorized subsets** of players can retrieve it if they **collaborate**

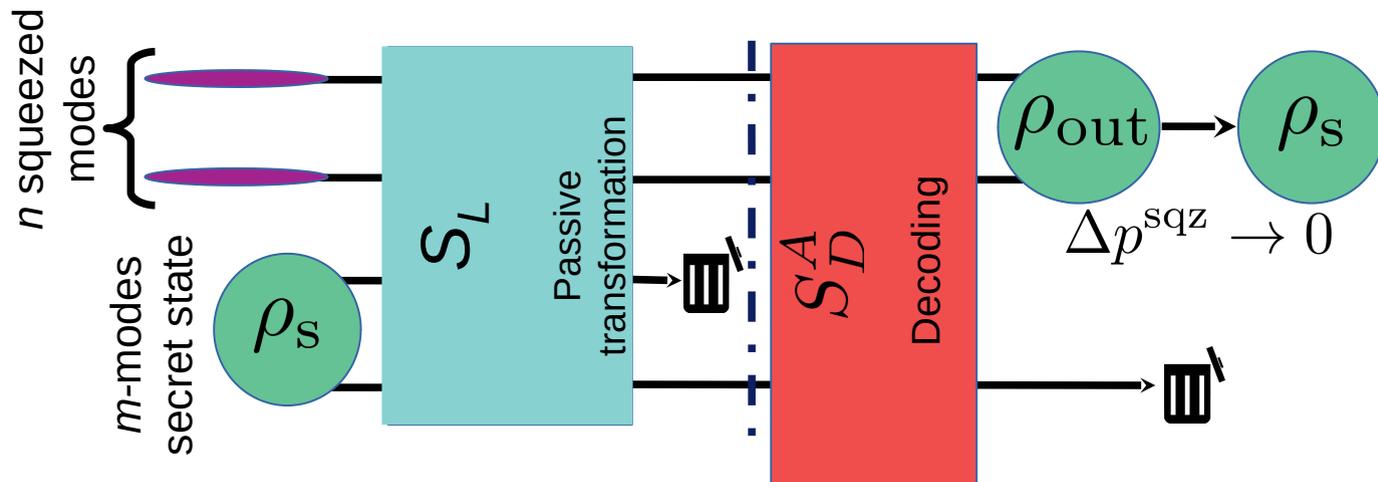


QQ: The secret is a quantum state

How it started: NatComm 8, 15645 (2017)

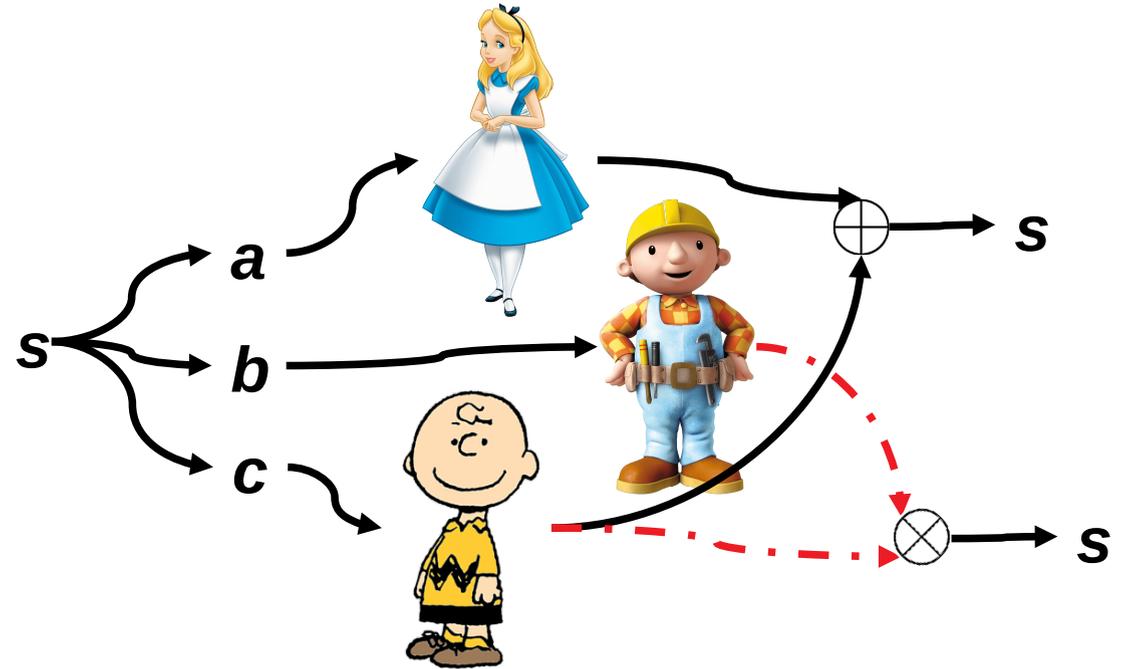
How it continued: PRA 100, 022303 (2019)

General CV (Gaussian) scheme



CV quantum state sharing: random codes

A **dealer** shares a **secret** with several **players** such that only **authorized subsets** of players can retrieve it if they **collaborate**

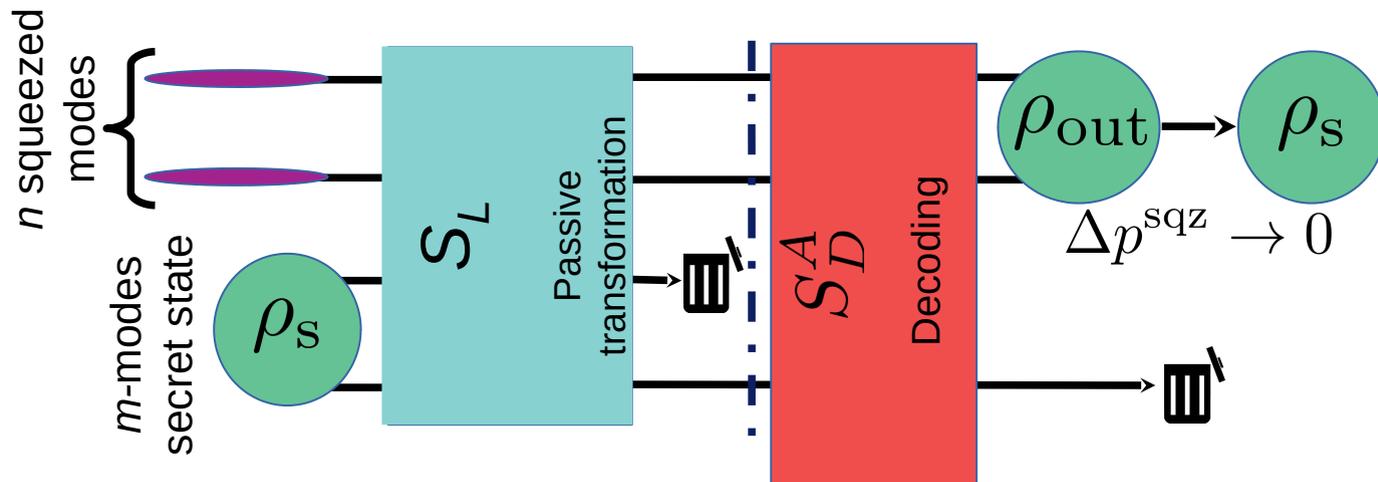


QQ: The secret is a quantum state

How it started: NatComm 8, 15645 (2017)

How it continued: PRA 100, 022303 (2019)

General CV (Gaussian) scheme



Our contribution:

A CV-QSS scheme can be realized by mixing the secret (quantum) state with squeezed states in **almost any** passive interferometer

- Generalizes previous protocols
- Experimentally friendly
- Analogous to erasure correcting

Current projects

(to appear soon)

With

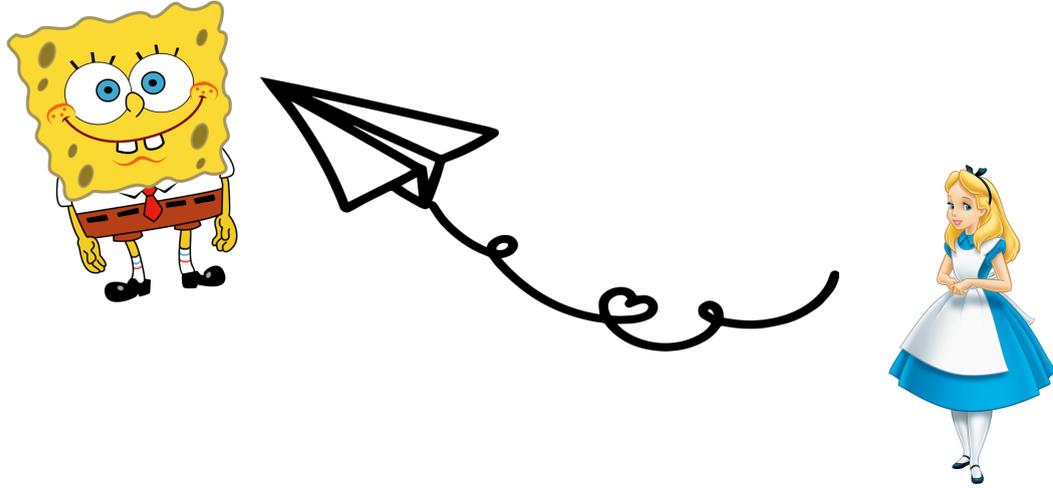


Jens Eisert



Jonathan Conrad

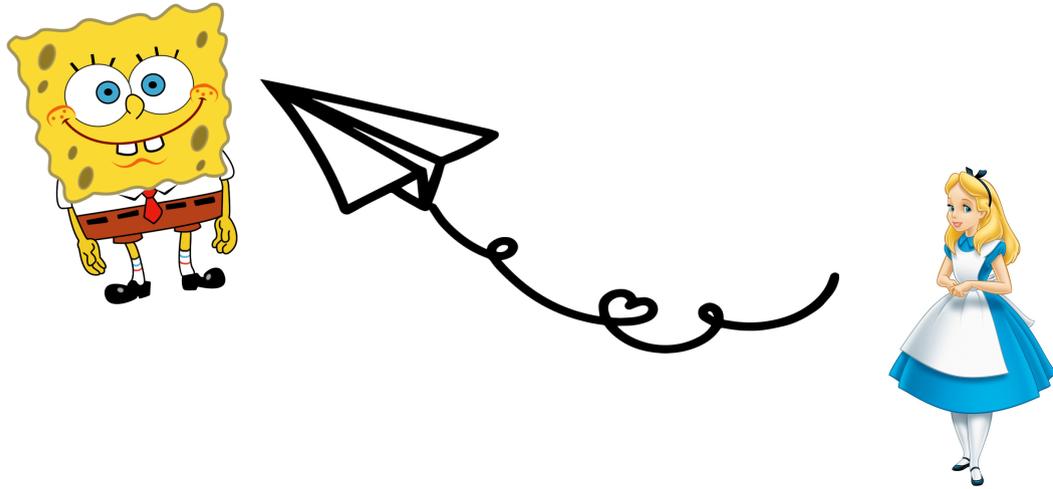
(Quantum) Error correction and harmonic oscillators



Information is always encoded in phys. syst.

➡ Always subject to **noise**

(Quantum) Error correction and harmonic oscillators

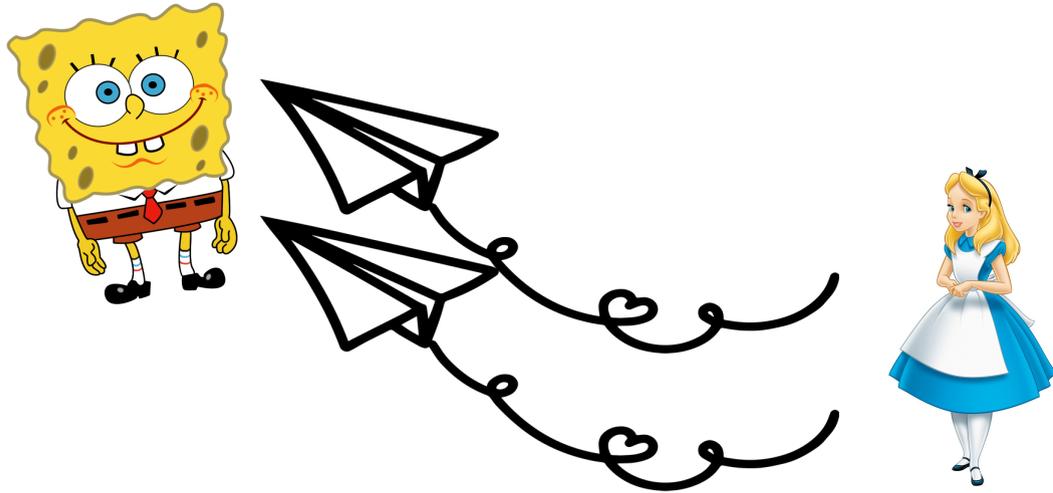


Information is always encoded in phys. syst.

➡ Always subject to **noise**

Rich theory of **error correction** to recover noisy messages:

(Quantum) Error correction and harmonic oscillators



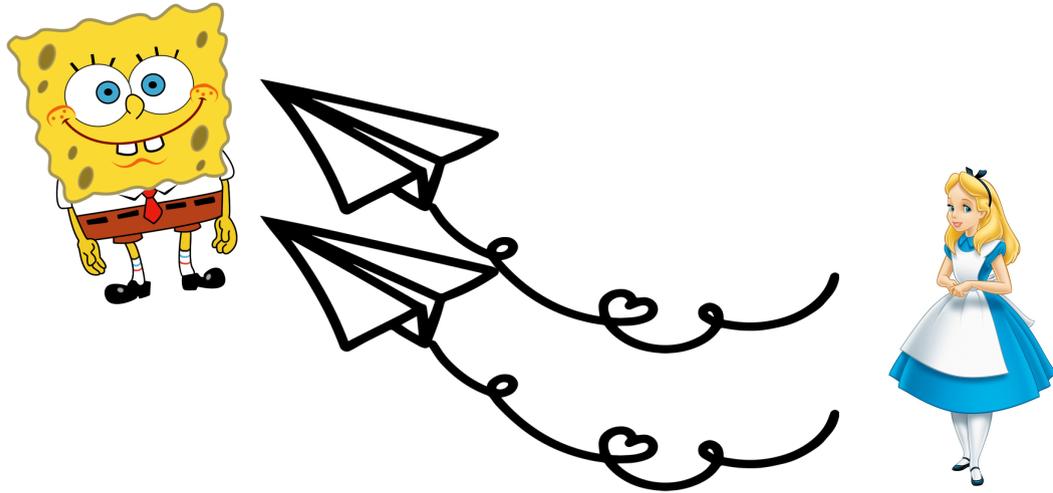
Information is always encoded in phys. syst.

➡ Always subject to **noise**

Rich theory of **error correction** to recover noisy messages:

Embed information in **larger system**

(Quantum) Error correction and harmonic oscillators



Information is always encoded in phys. syst.

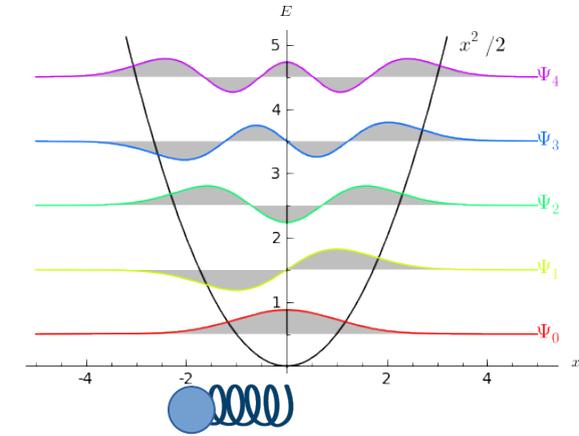
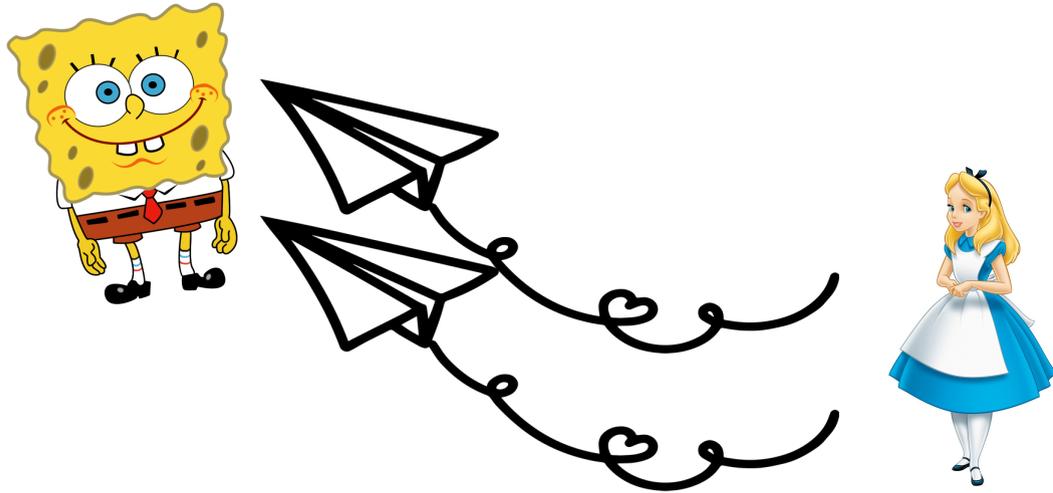
➡ Always subject to **noise**

Rich theory of **error correction** to recover noisy messages:

Embed information in **larger system**

Quantum: *mostly* qubits $\alpha |0\rangle + \beta |1\rangle$

(Quantum) Error correction and harmonic oscillators



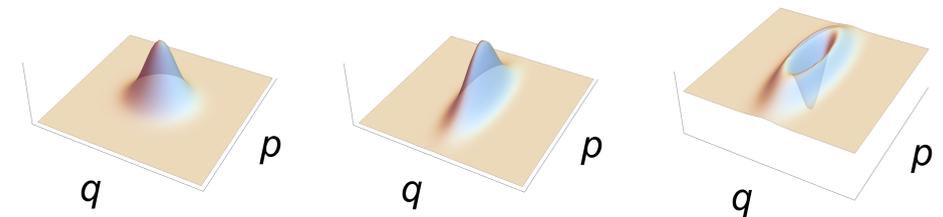
EM field mode, LC circuit, ...

Information is always encoded in phys. syst.

➡ Always subject to **noise**

In phase space:
Wigner Function

Rich theory of **error correction** to recover noisy messages:

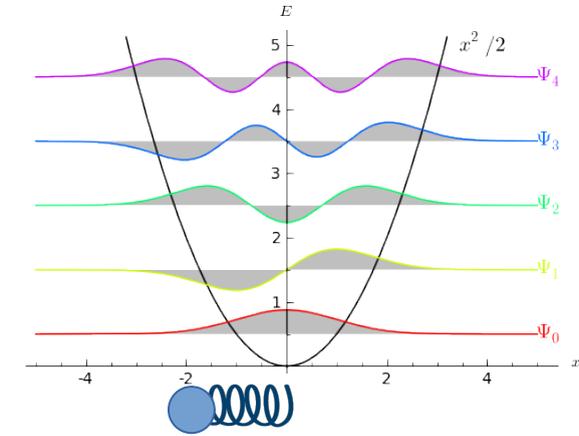
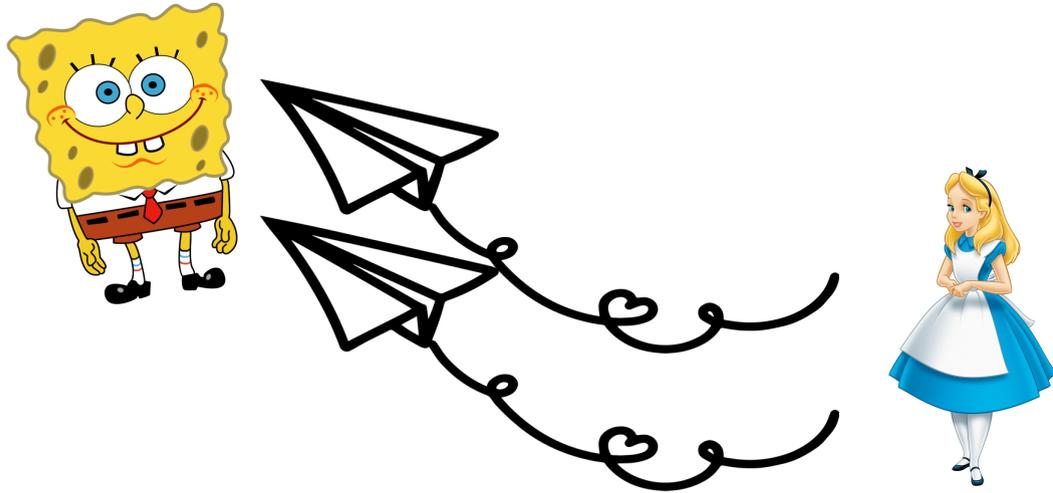


Quasi-probability distribution

Embed information in **larger system**

Quantum: *mostly* qubits $\alpha |0\rangle + \beta |1\rangle$

(Quantum) Error correction and harmonic oscillators



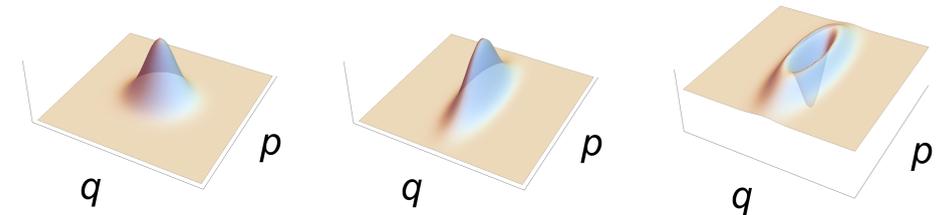
EM field mode, LC circuit, ...

Information is always encoded in phys. syst.

➡ Always subject to **noise**

In phase space:
Wigner Function

Rich theory of **error correction** to recover noisy messages:



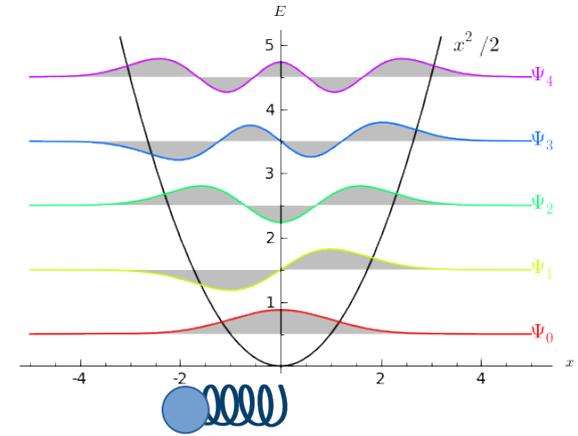
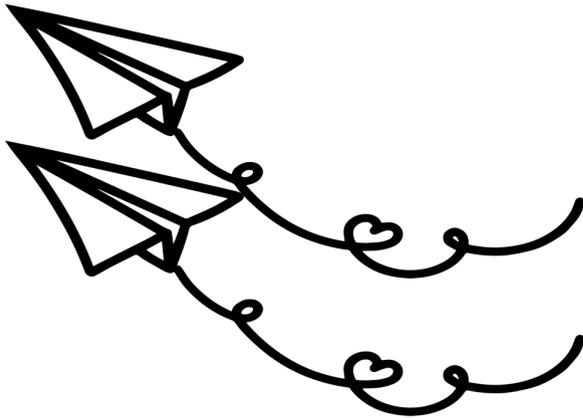
Embed information in **larger system**

Quasi-probability distribution

Quantum: *mostly* qubits $\alpha |0\rangle + \beta |1\rangle$

Infinitely many symbols!
How to restrict to finite?

(Quantum) Error correction and harmonic oscillators



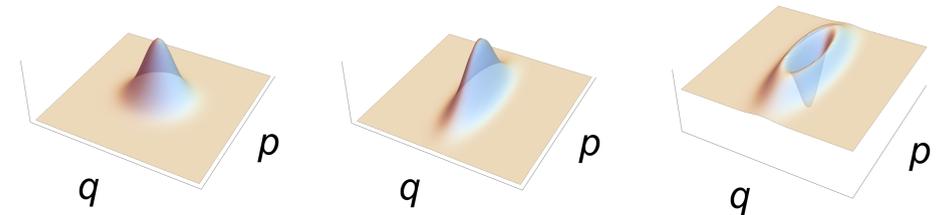
EM field mode, LC circuit, ...

Information is always encoded in phys. syst.

➡ Always subject to **noise**

In phase space:
Wigner Function

Rich theory of **error correction** to recover noisy messages:



Embed information in **larger system**

Quasi-probability distribution

Quantum: *mostly* qubits $\alpha |0\rangle + \beta |1\rangle$

Infinitely many symbols!
How to restrict to finite?

Symmetries!

Encoding qubits on a lattice

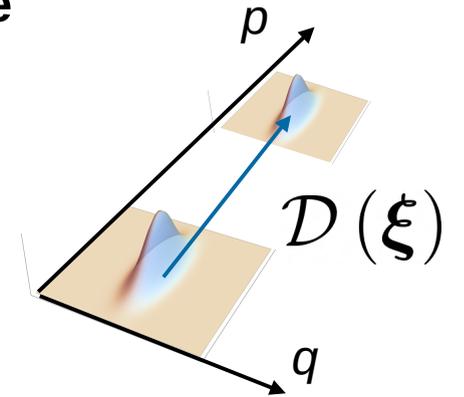
“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: **stabilized** by (commuting) **displacement** operators \rightarrow underlying **lattice**

Gottesman, Kitaev, Preskill PRA 64 (2001)



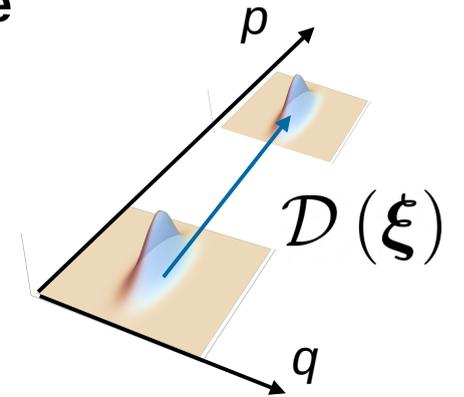
Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: stabilized by (commuting) displacement operators \rightarrow underlying lattice

Gottesman, Kitaev, Preskill PRA 64 (2001)

$$\mathcal{S} = \langle \mathcal{D}(\xi_1), \dots, \mathcal{D}(\xi_{2n}) \rangle \Rightarrow \text{Code: } \mathcal{D}(\xi_j) |\psi\rangle = |\psi\rangle$$



Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: stabilized by (commuting) displacement operators → underlying lattice

Gottesman, Kitaev, Preskill PRA 64 (2001)

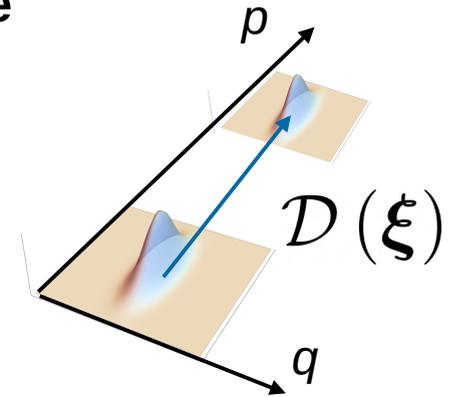
$$\mathcal{S} = \langle \mathcal{D}(\xi_1), \dots, \mathcal{D}(\xi_{2n}) \rangle \Rightarrow \text{Code: } \mathcal{D}(\xi_j) |\psi\rangle = |\psi\rangle$$

Logical states “live” on \mathcal{L}^*

$$\mathcal{D}(\xi_j) \mathcal{D}(\xi_k) = \mathcal{D}(\xi_j + \xi_k) (!)$$

→ Lattice of translations \mathcal{L}

→ Logical operations: dual \mathcal{L}^*



Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: stabilized by (commuting) displacement operators → underlying lattice

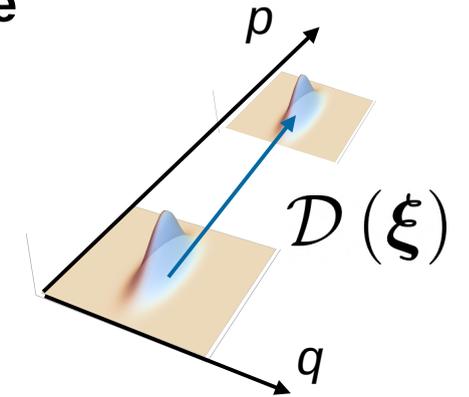
Gottesman, Kitaev, Preskill PRA 64 (2001)

$$\mathcal{S} = \langle \mathcal{D}(\xi_1), \dots, \mathcal{D}(\xi_{2n}) \rangle \Rightarrow \text{Code: } \mathcal{D}(\xi_j) |\psi\rangle = |\psi\rangle$$

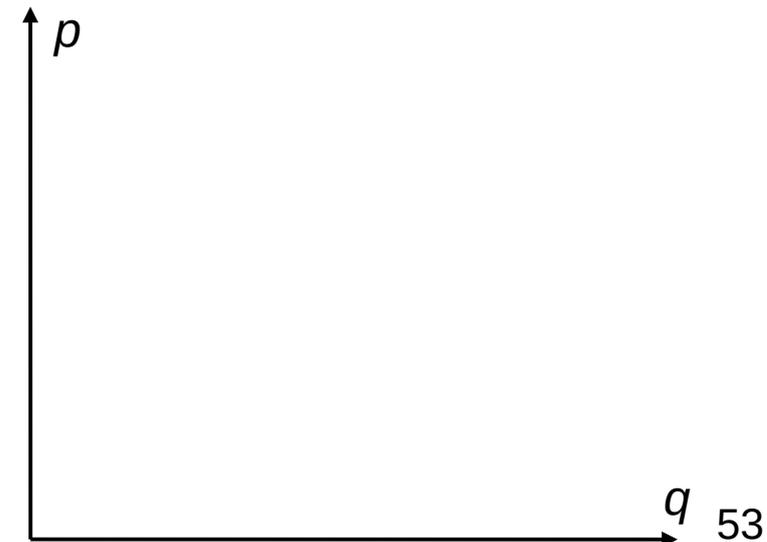
Logical states “live” on \mathcal{L}^*

$$\mathcal{D}(\xi_j) \mathcal{D}(\xi_k) = \mathcal{D}(\xi_j + \xi_k) (!)$$

- Lattice of translations \mathcal{L}
- Logical operations: dual \mathcal{L}^*



$$S_q = e^{i2\sqrt{\pi}\hat{q}} \quad S_p = e^{-i2\sqrt{\pi}\hat{p}}$$



Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: stabilized by (commuting) displacement operators → underlying lattice

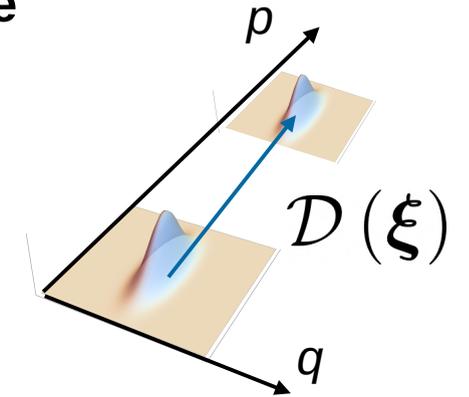
Gottesman, Kitaev, Preskill PRA 64 (2001)

$$\mathcal{S} = \langle \mathcal{D}(\xi_1), \dots, \mathcal{D}(\xi_{2n}) \rangle \Rightarrow \text{Code: } \mathcal{D}(\xi_j) |\psi\rangle = |\psi\rangle$$

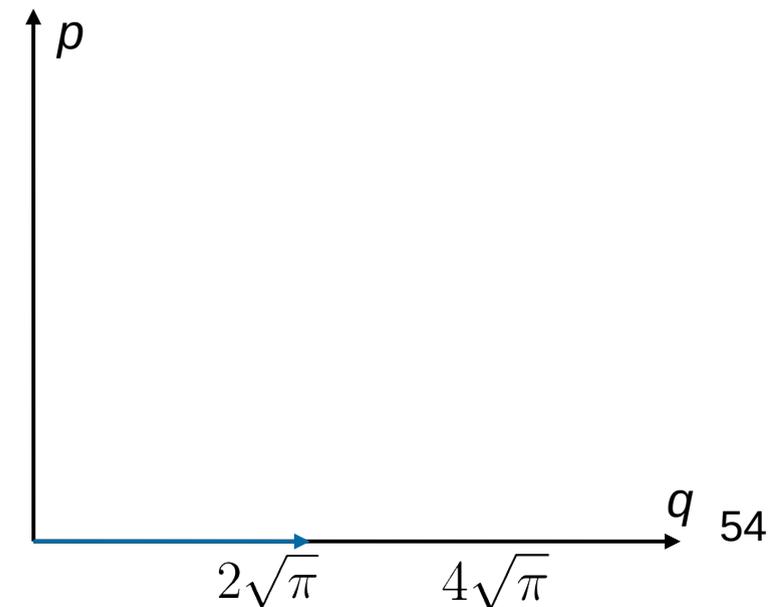
Logical states “live” on \mathcal{L}^*

$$\mathcal{D}(\xi_j) \mathcal{D}(\xi_k) = \mathcal{D}(\xi_j + \xi_k) (!)$$

- Lattice of translations \mathcal{L}
- Logical operations: dual \mathcal{L}^*



$$S_q = e^{i2\sqrt{\pi}\hat{q}} \quad S_p = e^{-i2\sqrt{\pi}\hat{p}}$$



Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: stabilized by (commuting) displacement operators → underlying lattice

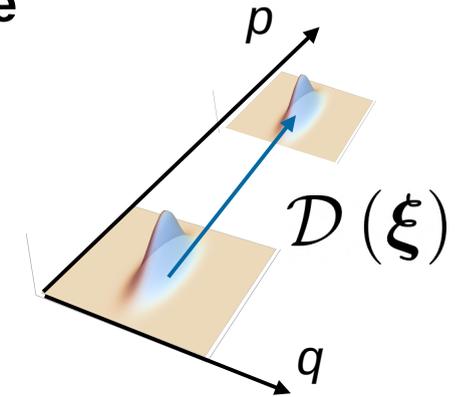
Gottesman, Kitaev, Preskill PRA 64 (2001)

$$\mathcal{S} = \langle \mathcal{D}(\xi_1), \dots, \mathcal{D}(\xi_{2n}) \rangle \Rightarrow \text{Code: } \mathcal{D}(\xi_j) |\psi\rangle = |\psi\rangle$$

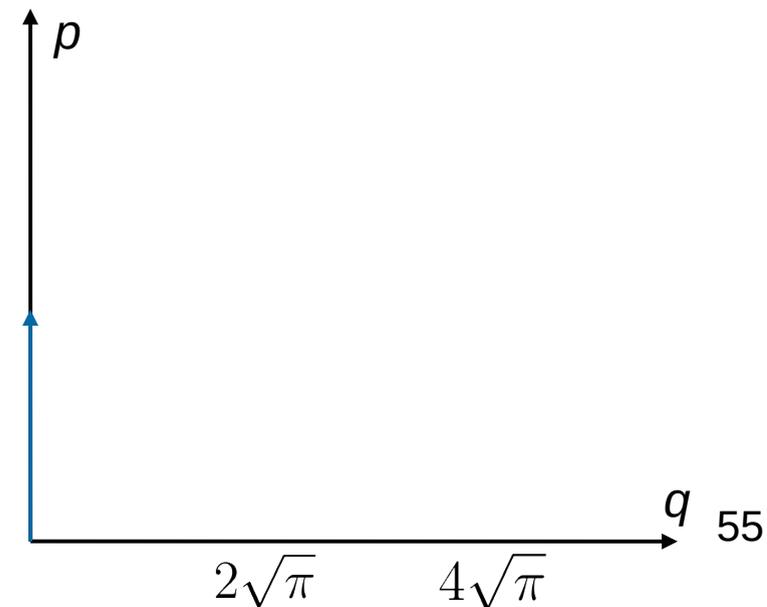
Logical states “live” on \mathcal{L}^*

$$\mathcal{D}(\xi_j) \mathcal{D}(\xi_k) = \mathcal{D}(\xi_j + \xi_k) (!)$$

- Lattice of translations \mathcal{L}
- Logical operations: dual \mathcal{L}^*



$$S_q = e^{i2\sqrt{\pi}\hat{q}} \quad S_p = e^{-i2\sqrt{\pi}\hat{p}}$$



Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: stabilized by (commuting) displacement operators → underlying lattice

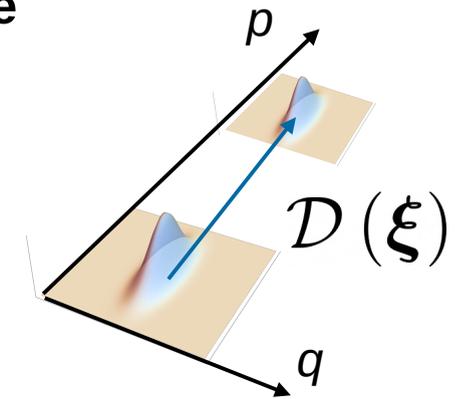
Gottesman, Kitaev, Preskill PRA 64 (2001)

$$\mathcal{S} = \langle \mathcal{D}(\xi_1), \dots, \mathcal{D}(\xi_{2n}) \rangle \Rightarrow \text{Code: } \mathcal{D}(\xi_j) |\psi\rangle = |\psi\rangle$$

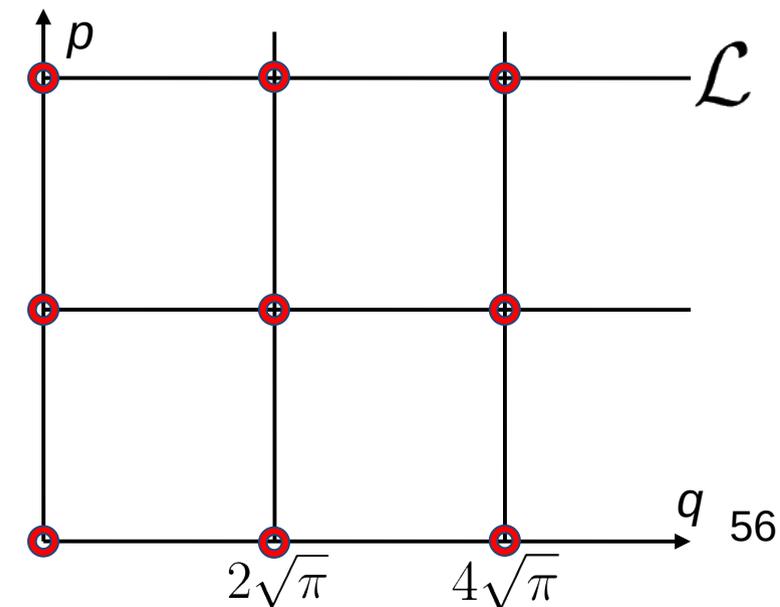
Logical states “live” on \mathcal{L}^*

$$\boxed{\mathcal{D}(\xi_j) \mathcal{D}(\xi_k) = \mathcal{D}(\xi_j + \xi_k)} \quad (!)$$

- Lattice of translations \mathcal{L}
- Logical operations: dual \mathcal{L}^*



$$S_q = e^{i2\sqrt{\pi}\hat{q}} \quad S_p = e^{-i2\sqrt{\pi}\hat{p}}$$



Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: stabilized by (commuting) displacement operators → underlying lattice

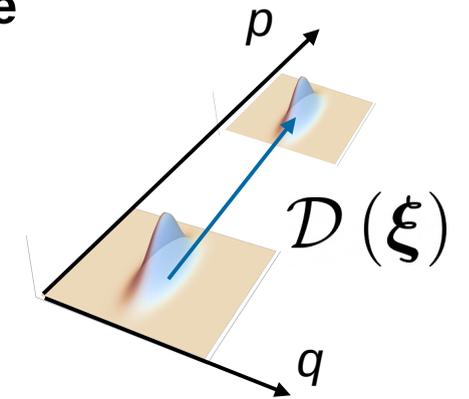
Gottesman, Kitaev, Preskill PRA 64 (2001)

$$\mathcal{S} = \langle \mathcal{D}(\xi_1), \dots, \mathcal{D}(\xi_{2n}) \rangle \Rightarrow \text{Code: } \mathcal{D}(\xi_j) |\psi\rangle = |\psi\rangle$$

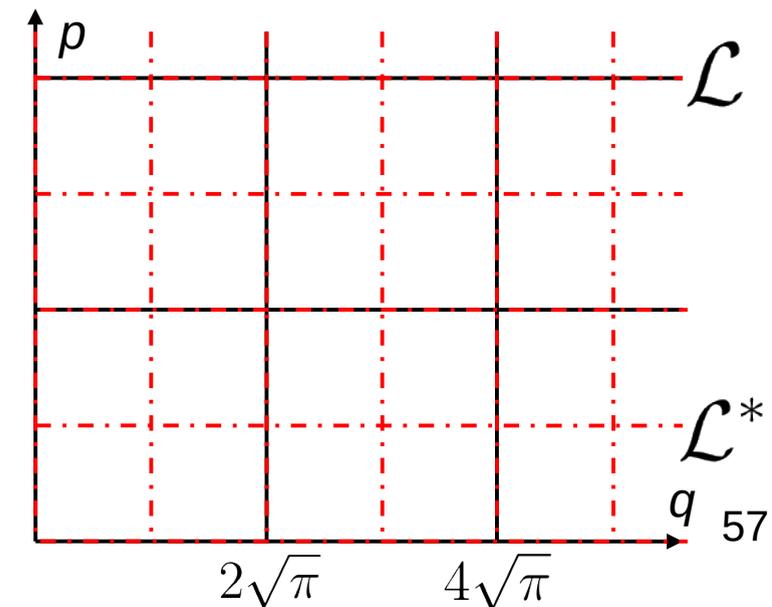
Logical states “live” on \mathcal{L}^*

$$\mathcal{D}(\xi_j) \mathcal{D}(\xi_k) = \mathcal{D}(\xi_j + \xi_k) \quad (!)$$

- Lattice of translations \mathcal{L}
- Logical operations: dual \mathcal{L}^*



$$S_q = e^{i2\sqrt{\pi}\hat{q}} \quad S_p = e^{-i2\sqrt{\pi}\hat{p}}$$



Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: stabilized by (commuting) displacement operators → underlying lattice

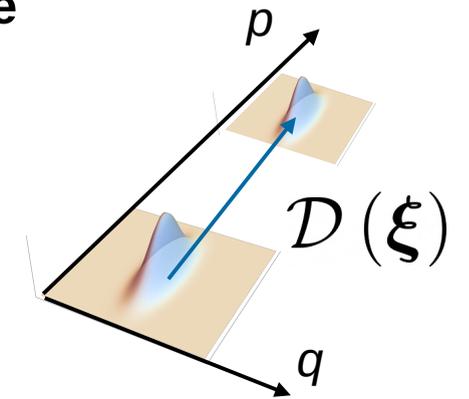
Gottesman, Kitaev, Preskill PRA 64 (2001)

$$\mathcal{S} = \langle \mathcal{D}(\xi_1), \dots, \mathcal{D}(\xi_{2n}) \rangle \Rightarrow \text{Code: } \mathcal{D}(\xi_j) |\psi\rangle = |\psi\rangle$$

Logical states “live” on \mathcal{L}^*

$$\mathcal{D}(\xi_j) \mathcal{D}(\xi_k) = \mathcal{D}(\xi_j + \xi_k) (!)$$

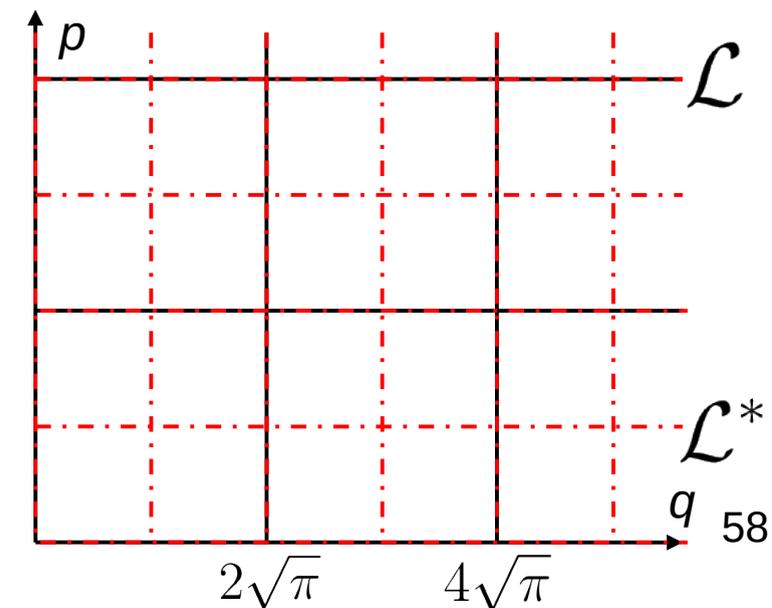
- Lattice of translations \mathcal{L}
- Logical operations: dual \mathcal{L}^*



$$S_q = e^{i2\sqrt{\pi}\hat{q}} \quad S_p = e^{-i2\sqrt{\pi}\hat{p}}$$

- Good protection against common noise processes

Albert et al, PRA 97 (2018)



Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: stabilized by (commuting) displacement operators → underlying lattice

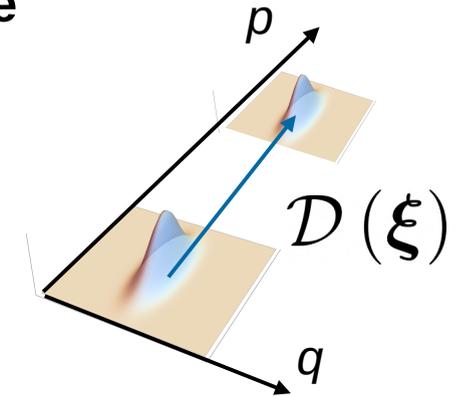
Gottesman, Kitaev, Preskill PRA 64 (2001)

$$\mathcal{S} = \langle \mathcal{D}(\xi_1), \dots, \mathcal{D}(\xi_{2n}) \rangle \Rightarrow \text{Code: } \mathcal{D}(\xi_j) |\psi\rangle = |\psi\rangle$$

Logical states “live” on \mathcal{L}^*

$$\mathcal{D}(\xi_j) \mathcal{D}(\xi_k) = \mathcal{D}(\xi_j + \xi_k) (!)$$

- Lattice of translations \mathcal{L}
- Logical operations: dual \mathcal{L}^*

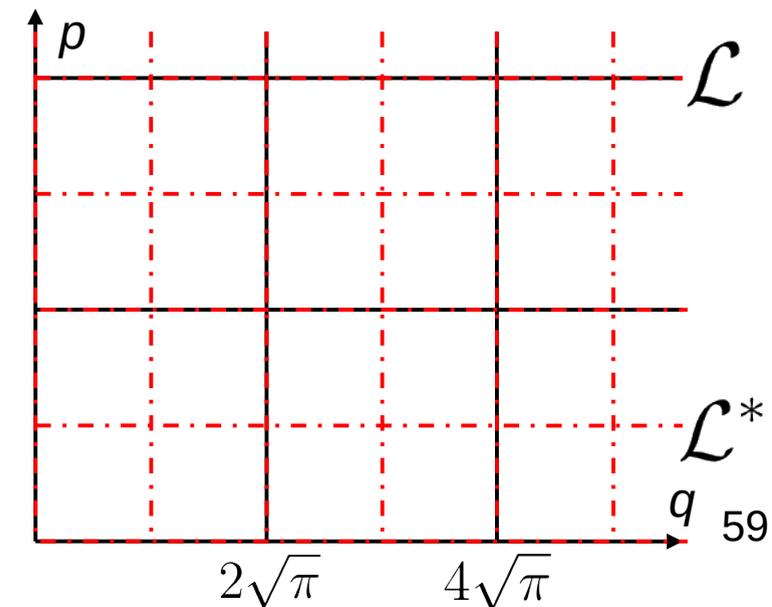


$$S_q = e^{i2\sqrt{\pi}\hat{q}} \quad S_p = e^{-i2\sqrt{\pi}\hat{p}}$$

- Good protection against common noise processes

Albert et al, PRA 97 (2018)

- Logical Clifford = Gaussian operations (“easy” - good for EC & QIP)



Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: stabilized by (commuting) displacement operators → underlying lattice

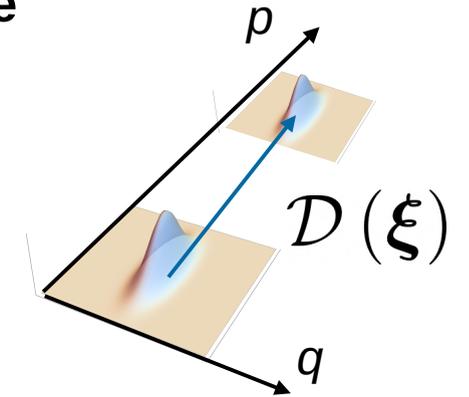
Gottesman, Kitaev, Preskill PRA 64 (2001)

$$\mathcal{S} = \langle \mathcal{D}(\xi_1), \dots, \mathcal{D}(\xi_{2n}) \rangle \implies \text{Code: } \mathcal{D}(\xi_j) |\psi\rangle = |\psi\rangle$$

Logical states “live” on \mathcal{L}^*

$$\mathcal{D}(\xi_j) \mathcal{D}(\xi_k) = \mathcal{D}(\xi_j + \xi_k) \quad (!)$$

- Lattice of translations \mathcal{L}
- Logical operations: dual \mathcal{L}^*



$$S_q = e^{i2\sqrt{\pi}\hat{q}} \quad S_p = e^{-i2\sqrt{\pi}\hat{p}}$$

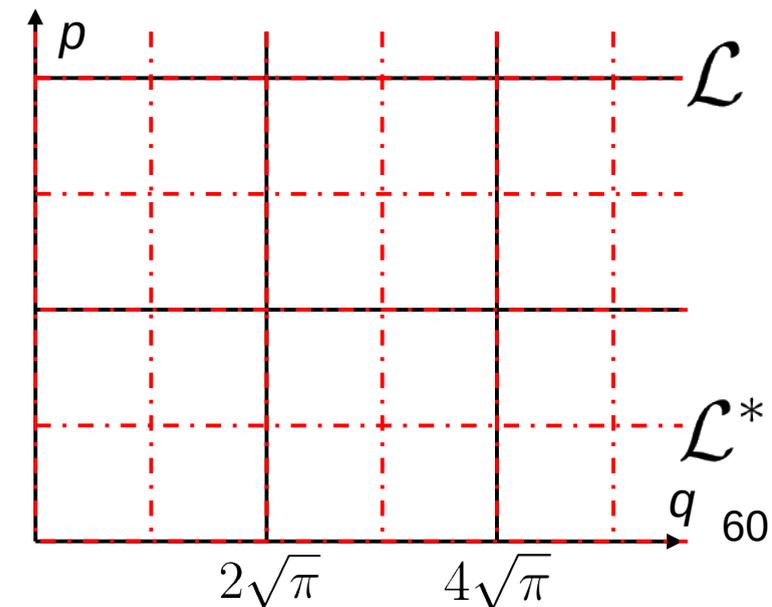
- Good protection against common noise processes

Albert et al, PRA 97 (2018)

- Logical Clifford = Gaussian operations (“easy” - good for EC & QIP)

- Can be used as effective qubits and combined with stabilizer codes

Vuillot et al, PRA 99 (2019) Noh&Chamberland PRA 101 (2020) Bourassa et al, Quantum 5 (2021)



Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: stabilized by (commuting) displacement operators → underlying lattice

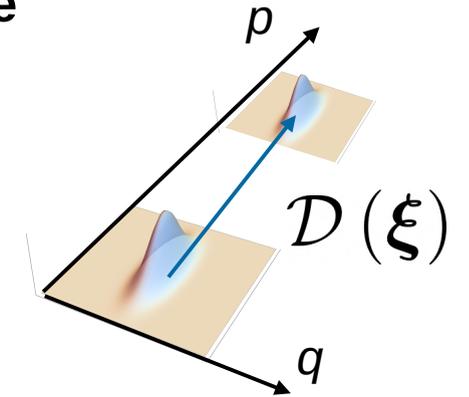
Gottesman, Kitaev, Preskill PRA 64 (2001)

$$\mathcal{S} = \langle \mathcal{D}(\xi_1), \dots, \mathcal{D}(\xi_{2n}) \rangle \Rightarrow \text{Code: } \mathcal{D}(\xi_j) |\psi\rangle = |\psi\rangle$$

Logical states “live” on \mathcal{L}^*

$$\mathcal{D}(\xi_j) \mathcal{D}(\xi_k) = \mathcal{D}(\xi_j + \xi_k) (!)$$

- Lattice of translations \mathcal{L}
- Logical operations: dual \mathcal{L}^*



$$S_q = e^{i2\sqrt{\pi}\hat{q}} \quad S_p = e^{-i2\sqrt{\pi}\hat{p}}$$

- Good protection against common noise processes

Albert et al, PRA 97 (2018)

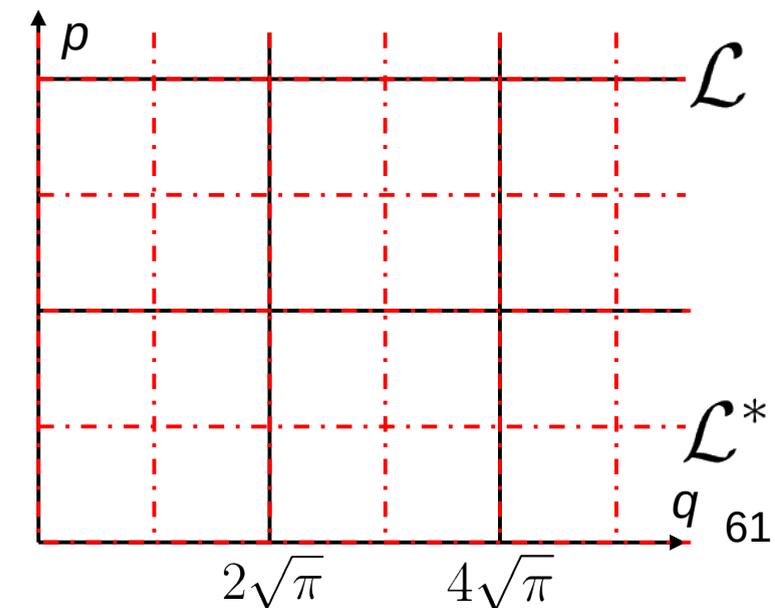
- Logical Clifford = Gaussian operations (“easy” - good for EC & QIP)

- Can be used as effective qubits and combined with stabilizer codes

Vuillot et al, PRA 99 (2019) Noh&Chamberland PRA 101 (2020) Bourassa et al, Quantum 5 (2021)

- Can protect CV systems (idea: error mitigation for Boson Sampling)

Noh et al, PRL 125 (2020)



Encoding qubits on a lattice

“Logical subspace”: finite photon number, odd/even photon number, or...**translation symmetries!**

Grid codes: stabilized by (commuting) displacement operators → underlying lattice

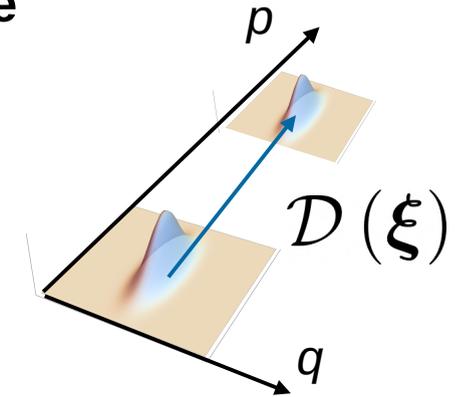
Gottesman, Kitaev, Preskill PRA 64 (2001)

$$\mathcal{S} = \langle \mathcal{D}(\xi_1), \dots, \mathcal{D}(\xi_{2n}) \rangle \Rightarrow \text{Code: } \mathcal{D}(\xi_j) |\psi\rangle = |\psi\rangle$$

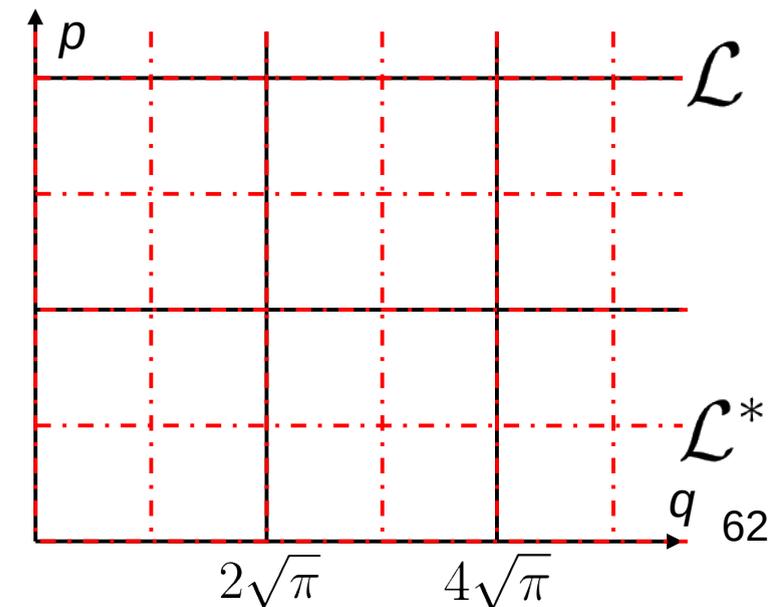
Logical states “live” on \mathcal{L}^*

$$\mathcal{D}(\xi_j) \mathcal{D}(\xi_k) = \mathcal{D}(\xi_j + \xi_k) (!)$$

- Lattice of translations \mathcal{L}
- Logical operations: dual \mathcal{L}^*



$$S_q = e^{i2\sqrt{\pi}\hat{q}} \quad S_p = e^{-i2\sqrt{\pi}\hat{p}}$$



- Good protection against common noise processes
Albert et al, PRA 97 (2018)
- Logical Clifford = Gaussian operations (“easy” - good for EC & QIP)
- Can be used as effective qubits and combined with stabilizer codes
Vuillot et al, PRA 99 (2019) Noh&Chamberland PRA 101 (2020) Bourassa et al, Quantum 5 (2021)
- Can protect CV systems (idea: error mitigation for Boson Sampling)
Noh et al, PRL 125 (2020)
- Logical states thought hard to realize, now there are experiments!
Flühmann et al, Nature 566 (2019) Campagne-Ibarcq et al, Nature 584 (2020)

Leveraging the lattice point of view

Grid states are somewhat resistant to noise, but still need to add redundancy

Leveraging the lattice point of view

Grid states are somewhat resistant to noise, but still need to add redundancy

Up to now: *concatenation* → add “qubit level symmetries” over many grid-encoded oscillators
→ “lattice picture” only used for individual oscillators

Leveraging the lattice point of view

Grid states are somewhat resistant to noise, but still need to add redundancy

Up to now: *concatenation* → add “qubit level symmetries” over many grid-encoded oscillators
→ “lattice picture” only used for individual oscillators

Q: Can lattice properties be exploited more?

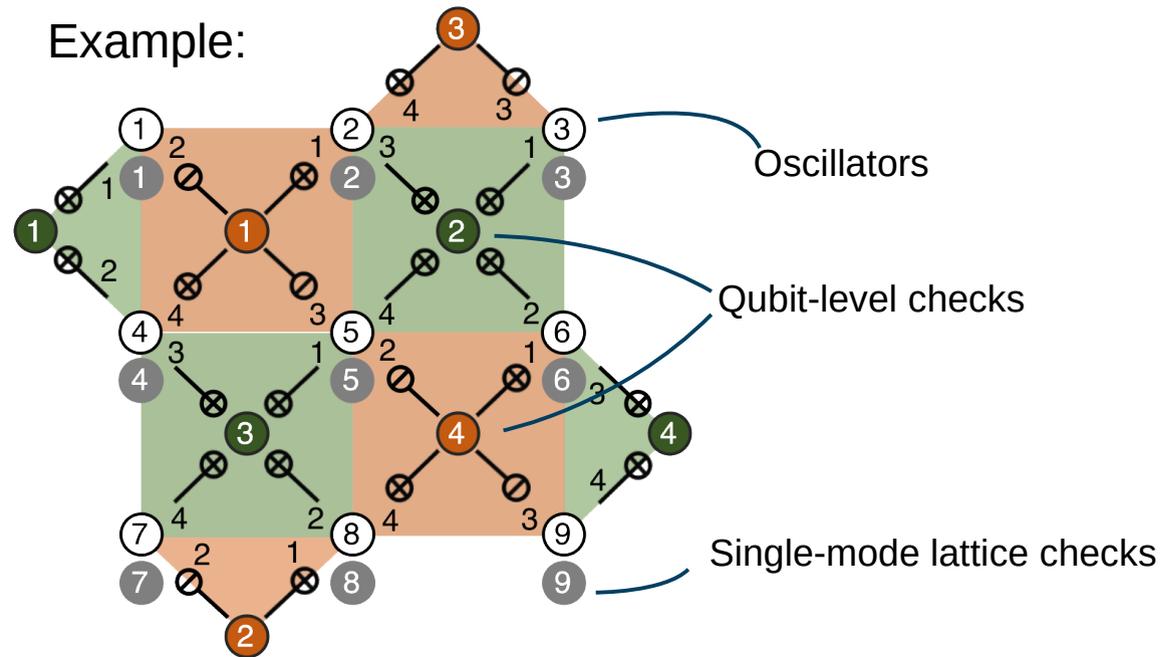
- Rich theory of lattices
- Advantages:
 - Better resource use
 - New codes
 - Additional proof techniques
 - Better decoding techniques?

Leveraging the lattice point of view

Grid states are somewhat resistant to noise, but still need to add redundancy

Up to now: *concatenation* → add “qubit level symmetries” over many grid-encoded oscillators
 → “lattice picture” only used for individual oscillators

Example:



Noh&Chamberland PRA 101 (2020)

$2n = 18$ mode-wise stab.

+

$n-1 = 8$ qubit stab.

 Tot = $3n - 1 = 26$ stab.



Can achieve same code with **18** total stab.

Q: Can lattice properties be exploited more?

- Rich theory of lattices
- Advantages:
 - Better resource use
 - New codes
 - Additional proof techniques
 - Better decoding techniques?