Gottesman-Kitaev-Preskill bosonic error correcting codes: a lattice perspective

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Jonathan Conrad, Jens Eisert, Francesco Arzani













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Error correction: redundancy



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Quantum: mostly qubits $\alpha \left| 0 \right\rangle + \beta \left| 1 \right\rangle$



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Quantum: *mostly* qubits $\alpha |0\rangle + \beta |1\rangle$ **Bosonic codes**: *oscillators*



EM field mode, LC circuit, ...



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ΕM

$$\hat{\boldsymbol{x}} = (q_1, \dots, q_n, p_1, \dots, p_n)^T$$

$$[\hat{\boldsymbol{x}}_j, \hat{\boldsymbol{x}}_k] = iJ_{jk}$$
field mode, LC circuit, ...
$$\hat{\boldsymbol{x}} = \begin{pmatrix} \mathbf{0} & \mathbb{I} \\ -\mathbb{I} & \mathbf{0} \end{pmatrix}$$



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$$J = \begin{pmatrix} \mathbf{0} & \mathbb{I} \\ -\mathbb{I} & \mathbf{0} \end{pmatrix}$$

Phase space:

$$W_{\rho}(\boldsymbol{q},\boldsymbol{p}) = (2\pi)^{-2n} \int \mathrm{d}^{n}\boldsymbol{y} \left\langle \boldsymbol{q} - \frac{\boldsymbol{y}}{2} \right| \hat{\rho} \left| \boldsymbol{q} + \frac{\boldsymbol{y}}{2} \right\rangle e^{i\boldsymbol{p}\cdot\boldsymbol{y}}$$

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Information always encoded in phys. syst. Always subject to noise

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$$p_{q} \qquad p_{q} \qquad p_{q}$$



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Displacements:

$$D^{\dagger}(\boldsymbol{\xi}) \hat{\boldsymbol{x}} D(\boldsymbol{\xi}) = \hat{\boldsymbol{x}} + \sqrt{2\pi} \boldsymbol{\xi}$$

$$D^{\dagger}(\boldsymbol{\xi}) D(\boldsymbol{\xi}) = e^{-i2\pi \boldsymbol{\xi}^{T} J \boldsymbol{\eta}} D(\boldsymbol{\eta}) D(\boldsymbol{\xi})$$

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$$\mathcal{S}=\langle D\left(m{\xi}_{1}
ight),\ldots,D\left(m{\xi}_{2n}
ight)
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 Code: $D\left(m{\xi}_{j}
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ight
angle=\left|\psi
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 $\boldsymbol{\xi}_{j}^{T}J\boldsymbol{\xi}_{k}\in\mathbb{Z}$



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$$\mathcal{L}$$

→ Logical operations: dual \mathcal{L}^{\perp}



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$$q$$

$$\frac{q}{2\sqrt{\pi}} \qquad 4\sqrt{\pi}$$

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$$q \quad q$$

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- Good against common noise Albert et al, PRA 97 (2018)







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 - Can protect CV systems Noh et al, PRL 125 (2020)







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 - Can protect CV systems Noh et al, PRL 125 (2020)
 - Logical states experimentally accessible Flühmann et al, Nature 566 (2019) Campagne-Ibarcg et al, Nature 584 (2020)







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Q: Can lattice properties be exploited more?

upshot: lattices are very well studied!

Good lattice intros:

J. Conway and N. Sloane.

Sphere packings, lattices and groups, volume 290. 1988

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Cse 206a: Lattice algorithms and applications, 2014

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lattices are very well studied!

...but not so much for GKP!

Gottesman, Kitaev, Preskill PRA 64 (2001) Harrington, Preskill PRA 64 (2001) Hänggli, Heinze, König, PRA 102 (2020) Hänggli, König, arXiv:2102.05545 (2021)

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Outline

- 1. Lattice formalism
- 2. Code properties from lattice bases
- 3. Symplectic operations
- 4. Distance bounds for GKP codes
- 5. Decoding problem and Θ functions
- 6. GKP codes beyond concatenation

Lattice formalism

$$\mathcal{S} = \langle D(\boldsymbol{\xi}_1), \dots, D(\boldsymbol{\xi}_{2n}) \rangle$$
$$M = (\boldsymbol{\xi}_1, \dots, \boldsymbol{\xi}_{2n})^T$$

Lattice formalism

 $\mathcal{S} = \langle D(\boldsymbol{\xi}_1), \dots, D(\boldsymbol{\xi}_{2n}) \rangle \qquad A_{jk} = (MJM^T)_{jk} \in \mathbb{Z} \Rightarrow [D(\boldsymbol{\xi}_j), D(\boldsymbol{\xi}_k)] = 0$ $M = (\boldsymbol{\xi}_1, \dots, \boldsymbol{\xi}_{2n})^T$

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$$D(\boldsymbol{\xi})D(\boldsymbol{\eta}) = e^{-i\pi\boldsymbol{\xi}^T J\boldsymbol{\eta}} D\left(\boldsymbol{\xi} + \boldsymbol{\eta}\right)$$
$$D(\boldsymbol{\xi})D(\boldsymbol{\eta}) = e^{-i\pi\boldsymbol{\xi}^T J\boldsymbol{\eta}} D\left(\boldsymbol{\xi} + \boldsymbol{\eta}\right) \Longrightarrow (-1)^{f(\boldsymbol{a},M)} D\left(\boldsymbol{a}^T M\right) \in \mathcal{S} \ \forall \boldsymbol{a} \in \mathbb{Z}^{2n}$$

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$$\mathcal{S} \simeq \mathcal{L} = \left\{ \boldsymbol{\xi} \in \mathbb{R}^{2n} | \boldsymbol{\xi}^T = \boldsymbol{a}^T M, \ \boldsymbol{a} \in \mathbb{Z}^{2n}
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Change of basis: $M \mapsto UM \Rightarrow M^{\perp} \mapsto U^{-T}M^{\perp}$ \longrightarrow $A \mapsto \begin{pmatrix} 0 & D \\ -D & 0 \end{pmatrix}_{43}$ Gottesman, Kitaev, Preskill PRA 64 (2001)

Results

Exploit basis manipulations/properties to study codes

$$M = \left(\boldsymbol{\xi}_1, \dots, \boldsymbol{\xi}_{2n}\right)^T$$

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Resource savings from lattice basis reduction

Ex: L=3 surface code



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Example of inequivalent codes: 2 qubits in 2 modes can have $D = \operatorname{diag}(4, 1)$ or $D = \operatorname{diag}(2, 2)$

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Lattice Θ function:

$$\Theta_{\mathcal{L}}(z) = \sum_{\boldsymbol{x} \in \mathcal{L}} q^{\boldsymbol{x}^T \boldsymbol{x}} = \sum_{\delta \in \mathcal{D}} N_{\delta} q^{\delta}$$

$$egin{aligned} q &= e^{i\pi z} \ \mathcal{D} &= ig\{ \|m{x}\|_2^2,\,m{x}\in\mathcal{L}ig\} \end{aligned}$$



consider
$$Q_{\mathcal{L}}(z) := \Theta_{\mathcal{L}^{\perp}}(z) - \Theta_{\mathcal{L}}(z) = N_{\Delta^2} q^{\Delta^2} + \dots$$



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Distance of **qubit** code completely specified by weight distribution $\{A_i\}_{i=0}^{2n}$ of stabs.

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Distance of GKP code $\mathcal{L}(M)$ completely specified by distance distribution $(\mathcal{D}, N_{\delta})$

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Distance of **qubit** code completely specified by weight distribution $\{A_i\}_{i=0}^{2n}$ of stabs.

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upshot:

Θ functions are ubiquitous in lattice theory, estimate those for approx MLD!

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Consider concatenated codes

$$\mathcal{L}_{\text{conc}} = \bigoplus_{j=1}^{n} \mathcal{L}_{\text{loc},j} + \operatorname{span}_{\mathbb{Z}} G, \quad G \subset \bigoplus_{j} \mathcal{L}_{\text{loc},j}^{\perp}$$

With *G* sympl. integral.

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Conjecture: distance computed similarly to concatenated codes (but strongly depends on G)

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Theorem: distance bound

$$\max\left\{\frac{\Delta_1}{\lambda_n(\mathcal{L}_2)}, \frac{\Delta_2}{\lambda_2(\mathcal{L}_1)}\right\} \le \Delta_{\otimes} \le \Delta_1 \Delta_2$$

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Conclusions

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